

A Simple Proof of the GHS and Further Inequalities

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Abstract. We formulate and prove a general set of correlation inequalities for spin $-1/2$ Ising ferromagnets with pair interactions. One of these is the Griffiths-Hurst-Sherman inequality. The proof is obtained using Gaussian random variables.

1. Introduction

We consider a system of N Ising spins with ferromagnetic pair interactions and non-negative external magnetic field. The probability $\mu(\sigma)$ of any configuration $\sigma = (\sigma_1, \dots, \sigma_N)$, $\sigma_i = \pm 1$, is given by the formula $\mu(\sigma) = Z^{-1} \exp(-\beta H(\sigma))$, where $\beta = (kT)^{-1}$,

$$H(\sigma) = -\frac{1}{2} \sum_{i \neq j} J_{ij} \sigma_i \sigma_j - h \sum_i \sigma_i, \quad J_{ij} = J_{ji} \geq 0, \quad h \geq 0, \quad (1.1)$$

$$Z = \sum_{\{\sigma\}} \exp(-\beta H(\sigma)). \quad (1.2)$$

In the sequel, we set $\beta = 1$. Given spin sites i, j, k , we define the third Ursell function

$$u_3(i, j, k) \equiv \langle \sigma_i \sigma_j \sigma_k \rangle - \langle \sigma_i \rangle \langle \sigma_j \sigma_k \rangle - \langle \sigma_j \rangle \langle \sigma_i \sigma_k \rangle - \langle \sigma_k \rangle \langle \sigma_i \sigma_j \rangle + 2 \langle \sigma_i \rangle \langle \sigma_j \rangle \langle \sigma_k \rangle, \quad (1.3)$$

where the bracket $\langle \rangle$ denotes the expected value with respect to the measure μ .

The Griffiths-Hurst-Sherman inequality (hereafter GHS inequality) states that

$$u_3(i, j, k) \leq 0. \quad (1.4)$$

An important consequence of this inequality is that the average magnetization per site is a concave function of magnetic field h , a fact needed for the proof of certain critical point exponent inequalities [1]. It has also been used by Preston [2] to show the absence of phase transitions in the thermodynamic limit for $h \neq 0$.

Inequality (1.4) was first proved by Griffiths, Hurst, and Sherman [1] and later by Lebowitz [3]. Our proof is completely self-contained and, we believe, is much simpler. It is based on ideas introduced by Monroe and Siegert [4], who obtained simple proofs of the GKS inequalities [5]. Similar methods have also been used by Monroe [6] to prove certain FKG inequalities [7]. At the end of the next section, we mention additional new inequalities which are proved by the same technique.

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2. Sketch of Proof and Further Results

For convenience, we assume $i = 1, j = 2, k = 3$. Our proof is based on the identity [8]

$$\exp \left[\frac{1}{2} \sum_{i,j} \xi_i v_{ij} \xi_j \right] = (2\pi)^{-N/2} (\det v)^{-1/2} \int \dots \int \prod_{i=1}^N dx_i \exp \left[-\frac{1}{2} \sum_{i,j} x_i (v^{-1})_{ij} x_j + \sum_i \xi_i x_i \right], \quad (2.1)$$

valid for any symmetric, real, positive definite matrix v and for any N complex variables ξ_i . The right-hand side of (2.1) can be considered as the expected value $E_x[\exp \sum_{1 \leq i \leq N} x_i \xi_i]$ with respect to the Gaussian density function

$$f_v(\bar{x}) = (2\pi)^{-N/2} (\det v)^{-1/2} \exp \left[-\frac{1}{2} \sum_{i,j} x_i (v^{-1})_{ij} x_j \right], \quad (2.2)$$

where $\bar{x} = (x_1, \dots, x_N)$. If v is a non-negative matrix, then given integers $n_i \geq 0$, one can show

$$E_x \left[\prod_{i=1}^N (x_i)^{n_i} \right] \begin{cases} = 0 & \text{if } \sum_{i=1}^N n_i \text{ is odd,} \\ \geq 0 & \text{if } \sum_{i=1}^N n_i \text{ is even.} \end{cases} \quad (2.3)$$

To use (2.1), we identify the variables ξ_i with the spin variables σ_i and form a matrix $v = J$ with off-diagonal elements J_{ij} , diagonal elements all equal to a number $J_0 \equiv J_{ii}$ large enough to guarantee that J is positive definite. We then let (x_1, \dots, x_N) , (y_1, \dots, y_N) , (z_1, \dots, z_N) , (w_1, \dots, w_N) be independent sets of Gaussian random variables, where the random variables in each set have joint density function f_J . Writing \bar{E} to denote expectation with respect to the product measure $f_J(\bar{x}) f_J(\bar{y}) f_J(\bar{z}) f_J(\bar{w})$, we show that

$$u_3(1, 2, 3) = K \cdot \bar{E} \left[D \left(\frac{\partial}{\partial x_l}, \frac{\partial}{\partial y_l}, \frac{\partial}{\partial z_l} \right) g(\bar{x}, \bar{y}, \bar{z}, \bar{w}) \right] \quad (2.4)$$

where K is some positive constant, D denotes a sum of products of partial derivatives involving $\frac{\partial}{\partial x_l}, \frac{\partial}{\partial y_l}, \frac{\partial}{\partial z_l}, l = 1, 2, 3$, and g is a certain function of $\bar{x}, \bar{y}, \bar{z}, \bar{w}$. The crux of the proof is to reexpress (2.4) in terms of new variables $\bar{\alpha} = \{\alpha_l\}, \bar{\beta} = \{\beta_l\}, \bar{\gamma} = \{\gamma_l\},$ and $\bar{\delta} = \{\delta_l\}, l = 1, \dots, N$, obtained by a certain orthogonal transformation of $\bar{x}, \bar{y}, \bar{z}, \bar{w}$. A remarkable simplification then occurs, and we have

$$u_3(1, 2, 3) = 2K \cdot \bar{E} \left[\frac{\partial}{\partial \beta_1} \frac{\partial}{\partial \gamma_2} \frac{\partial}{\partial \delta_3} \tilde{g}(\alpha, \beta, \gamma, \delta) \right], \quad (2.5)$$

where \tilde{g} is the transformed g . The GHS inequality (1.4) then follows from (2.3).

Our method has the following consequence. Let \mathcal{D} denote a product of partial derivatives of the form

$$\prod_{i,j,k,l} \left(\frac{\partial}{\partial \alpha_i} \right)^{n_i^\alpha} \left(\frac{\partial}{\partial \beta_j} \right)^{n_j^\beta} \left(\frac{\partial}{\partial \gamma_k} \right)^{n_k^\gamma} \left(\frac{\partial}{\partial \delta_l} \right)^{n_l^\delta}, \quad (2.6)$$

where i, j, k, l range over all lattice points $\{1, \dots, N\}$ and $n_i^\alpha, n_j^\beta, n_k^\gamma, n_l^\delta$ are non-negative integers. Each such $\tilde{\mathcal{D}}$ gives rise to a sum $c_{\mathcal{D}}$ of correlations.

Theorem 1. *Define*

$$N_1 = \sum_{i=1}^N n_i^\alpha, \quad N_2 = \sum_{j=1}^N n_j^\beta, \quad N_3 = \sum_{k=1}^N n_k^\gamma, \quad \text{and} \quad N_4 = \sum_{l=1}^N n_l^\delta.$$

We then have the following:

- a) $c_{\mathcal{D}} \leq 0$, if each $N_i, i = 1, \dots, 4$, is odd;
 - b) $c_{\mathcal{D}} \geq 0$, if each $N_i, i = 1, \dots, 4$, is even;
 - c) $c_{\mathcal{D}} = 0$, in all other cases.
- (2.7)

When $h > 0$, the oddness or evenness of N_1 is not to be considered.

The simplest of these $\tilde{\mathcal{D}}$ correspond to elementary GKS inequalities; e.g., $\tilde{\mathcal{D}} = \partial/\partial\alpha_i$ gives rise to $\langle \sigma_i \rangle \geq 0$; $\tilde{\mathcal{D}} = (\partial/\partial\beta_j)(\partial/\partial\beta_j)$ gives rise to $\langle \sigma_i \sigma_j \rangle - \langle \sigma_i \rangle \langle \sigma_j \rangle \geq 0$ (see Appendix for calculations). In future work, we shall investigate the explicit forms, in terms of the spin variables σ_i , of the inequalities in Theorem 1 as well as generalizations of them.

3. Proof of GHS Inequality and Theorem 1

We first note that the Boltzmann factor $\exp(-H(\sigma))$ can be written as

$$\begin{aligned} \exp(-H(\sigma)) &= E_{\bar{x}} \left[\exp \left\{ \sum_{i=1}^N \left([x_i + h] \sigma_i - \frac{1}{2} J_0 \sigma_i^2 \right) \right\} \right] \\ &= \exp \left(-\frac{NJ_0}{2} \right) E_{\bar{x}} \left[\prod_{i=1}^N \exp \{ (x_i + h) \sigma_i \} \right], \end{aligned} \quad (3.1)$$

since each $\sigma_i^2 = 1$. Thus,

$$Z = 2^N \exp \left(-\frac{NJ_0}{2} \right) E_{\bar{x}} \left[\prod_{i=1}^N \cosh(x_i + h) \right], \quad (3.2)$$

and, for example

$$\begin{aligned} \langle \sigma_1 \rangle \langle \sigma_2 \sigma_3 \rangle &= \left[\frac{1}{Z} \sum_{\{\sigma\}} \sigma_1 \exp(-H(\sigma)) \right] \left[\frac{1}{Z} \sum_{\{\sigma'\}} \sigma'_2 \sigma'_3 \exp(-H(\sigma')) \right] \\ &= 2^{2N} Z^{-2} \exp(-NJ_0) \left\{ \sum_{\{\sigma\}} E_{\bar{x}} \left[\frac{\partial}{\partial x_1} \prod_{i=1}^N \exp[(x_i + h) \sigma_i] \right] \right. \\ &\quad \left. \sum_{\{\sigma'\}} E_{\bar{y}} \left[\frac{\partial}{\partial y_2} \frac{\partial}{\partial y_3} \prod_{i=1}^N \exp[(y_i + h) \sigma'_i] \right] \right\} \\ &= K \cdot \bar{E} \left[\frac{\partial}{\partial x_1} \frac{\partial}{\partial y_2} \frac{\partial}{\partial y_3} g(\bar{x}, \bar{y}, \bar{z}, \bar{w}) \right], \end{aligned} \quad (3.3)$$

where $K = 2^{4N} Z^{-4} \exp(-2NJ_0)$, and

$$g(\bar{x}, \bar{y}, \bar{z}, \bar{w}) = \prod_{i=1}^N \cosh(x_i + h) \cosh(y_i + h) \cosh(z_i + h) \cosh(w_i + h), \quad (3.4)$$

and where we have used (3.2) twice for the \bar{z} and \bar{w} variables. Therefore, to prove (2.4), we take the same K and g and define

$$D\left(\frac{\partial}{\partial x_i}, \frac{\partial}{\partial y_i}, \frac{\partial}{\partial z_i}\right) = \frac{\partial}{\partial x_1} \frac{\partial}{\partial x_2} \frac{\partial}{\partial x_3} - \frac{\partial}{\partial x_1} \frac{\partial}{\partial x_2} \frac{\partial}{\partial y_3} - \frac{\partial}{\partial x_1} \frac{\partial}{\partial y_2} \frac{\partial}{\partial y_3} \\ - \frac{\partial}{\partial x_1} \frac{\partial}{\partial y_2} \frac{\partial}{\partial x_3} + 2 \frac{\partial}{\partial x_1} \frac{\partial}{\partial y_2} \frac{\partial}{\partial z_3}. \quad (3.5)$$

The representation of $u_3(1, 2, 3)$ with this particular D is not unique, for D may be replaced by any differential operator obtained by performing the same permutation to the three sets of variables (x_i, y_i, z_i, w_i) , $i = 1, 2, 3$. This follows from the permutation invariance of $g(\bar{x}, \bar{y}, \bar{z}, \bar{w})$ and of $f_J(\bar{x})f_J(\bar{y})f_J(\bar{z})f_J(\bar{w})$.

For each $i = 1, \dots, N$, we define new variables $(\alpha_i, \beta_i, \gamma_i, \delta_i)$ by the formula

$$\begin{pmatrix} \alpha_i \\ \beta_i \\ \gamma_i \\ \delta_i \end{pmatrix} = \frac{1}{2} \begin{pmatrix} 1 & 1 & 1 & 1 \\ -1 & 1 & -1 & 1 \\ -1 & -1 & 1 & 1 \\ 1 & -1 & -1 & 1 \end{pmatrix} \begin{pmatrix} x_i \\ y_i \\ z_i \\ w_i \end{pmatrix}. \quad (3.6)$$

The orthogonal matrix in (3.6) is the direct product with itself of the matrix $2^{-\frac{1}{2}} \begin{pmatrix} 1 & 1 \\ -1 & 1 \end{pmatrix}$ which is analogous to the one used in [4]. In the Appendix we show that $2(\partial/\partial\beta_1)(\partial/\partial\gamma_2)(\partial/\partial\delta_3)$ corresponds to the differential operator D given by (3.5). Also, $g(\bar{x}, \bar{y}, \bar{z}, \bar{w})$ goes over to

$$\tilde{g}(\bar{\alpha}, \bar{\beta}, \bar{\gamma}, \bar{\delta}) = \prod_{i=1}^N t(\alpha_i, \beta_i, \gamma_i, \delta_i), \quad (3.7)$$

where

$$t(\alpha_i, \beta_i, \gamma_i, \delta_i) = \cosh 2(\alpha_i + 2h) + \cosh 2\beta_i + \cosh 2\gamma_i + \cosh 2\delta_i \\ + \cosh [-(\alpha_i + 2h) + \beta_i + \gamma_i + \delta_i] + \cosh [(\alpha_i + 2h) - \beta_i + \gamma_i + \delta_i] \\ + \cosh [(\alpha_i + 2h) + \beta_i - \gamma_i + \delta_i] + \cosh [(\alpha_i + 2h) + \beta_i + \gamma_i - \delta_i]. \quad (3.8)$$

Since $f_J(\bar{x})f_J(\bar{y})f_J(\bar{z})f_J(\bar{w})$ transforms to $f_J(\bar{\alpha})f_J(\bar{\beta})f_J(\bar{\gamma})f_J(\bar{\delta})$, (2.5) results. We claim that $t(\alpha_i, \beta_i, \gamma_i, \delta_i)$ is a sum of terms of the form

$$A(\alpha_i + 2h)^{j_\alpha} (\beta_i)^{j_\beta} (\gamma_i)^{j_\gamma} (\delta_i)^{j_\delta} - B(\alpha_i + 2h)^{k_\alpha} (\beta_i)^{k_\beta} (\gamma_i)^{k_\gamma} (\delta_i)^{k_\delta}, \quad (3.9)$$

where A and B are non-negative coefficients, $j_\alpha, j_\beta, j_\gamma, j_\delta$ are non-negative *even* integers, and $k_\alpha, k_\beta, k_\gamma, k_\delta$ are positive *odd* integers. This is proved at the end of this section. Given this, the GHS inequality (1.4) follows. Indeed, \tilde{g} is then a sum of terms of the form

$$C \prod_{i=1}^N (\alpha_i + 2h)^{l_i} (\beta_i)^{m_i} (\gamma_i)^{n_i} (\delta_i)^{p_i}, \quad (3.10)$$

where $\Sigma l_i, \Sigma m_i, \Sigma n_i$, and Σp_i are either all odd or all even and the constant C is either negative or positive, respectively. In the odd case, with $h \geq 0$, the derivative operator $(\partial/\partial\beta_1)(\partial/\partial\gamma_2)(\partial/\partial\delta_3)$ working on (3.10) either gives zero or converts

(3.10) into an expression containing an even number of β_i 's, γ_i 's, and δ_i 's. Also, when the term $\Pi_i(\alpha_i + 2h)^i$ is expanded, there appear an even number of α_i 's (as well as an odd number, which gives zero in (2.5)). Hence by (2.3) the contribution of (3.10) to the \bar{E} -expectation in (2.5) is non-positive. In the odd case with $h = 0$ and in the even case with any $h \geq 0$, the contribution of (3.10) to the \bar{E} -expectation in (2.5) is easily shown to be zero.

To prove Theorem 1, first assume $h > 0$. In cases a) and b), regardless of the parity of N_1 , $\hat{\mathcal{G}}$ operating on (3.10) either gives zero or converts (3.10) into an expression containing an even number of α_i 's, β_i 's, γ_i 's, and δ_i 's (again expand $\Pi_i(\alpha_i + 2h)^i$). Such a term gives a non-positive contribution in case a), a non-negative contribution in case b), to the \bar{E} -expectation in (2.5). In case c), there are no non-zero contributions. When $h = 0$, we must consider the parity of N_1 because then the α_i variables are on the same footing as the β_i , γ_i , and δ_i variables.

To prove (3.9), we expand the cosh terms in (3.8). To ease the notation, we drop the subscripts in α_i , β_i , γ_i , and δ_i . The first four terms in (3.8) give rise to terms of the form in (3.9), with $A > 0$ and $B = 0$. Each of the last four gives rise to terms of the form

$$A' [\varepsilon_1^i (\alpha + 2h) + \varepsilon_2^i \beta + \varepsilon_3^i \gamma + \varepsilon_4^i \delta]^{2n}, \quad (3.11)$$

$$A' \geq 0, n \geq 0, i = 1, 2, 3, 4,$$

where

$$\varepsilon_j^i = \begin{cases} -1, & j = i, \\ 1, & j \neq i. \end{cases} \quad (3.12)$$

But

$$\begin{aligned} & \sum_{i=1}^4 [\varepsilon_1^i (\alpha + 2h) + \varepsilon_2^i \beta + \varepsilon_3^i \gamma + \varepsilon_4^i \delta]^{2n} \\ &= B' \sum_{i=1}^4 \sum_{j=1}^4 [\varepsilon_1^i (\alpha + 2h)]^{m_\alpha} (\varepsilon_2^i \beta)^{m_\beta} (\varepsilon_3^i \gamma)^{m_\gamma} (\varepsilon_4^i \delta)^{m_\delta}, \end{aligned} \quad (3.13)$$

where $B' > 0$ and the outer summation extends over all non-negative integers $m_\alpha, m_\beta, m_\gamma, m_\delta$ with sum $2n$. Hence, either none, two, or all four of the m 's are odd. In the first and third cases, we have terms of the form in (3.9) ($A > 0, B = 0$, and $A = 0, B > 0$, respectively). If two of the m 's are odd, then when the inner sum over i is done, there results a zero. This completes the proof of the GHS inequality and of Theorem 1.

Appendix

We first do the calculation for the GHS inequalities, then for the GKS inequalities mentioned in the text. For convenience, rather than write partial derivatives like $(\partial/\partial\beta_1)(\partial/\partial\gamma_2)(\partial/\partial\delta_3)$ and $(\partial/\partial x_1)(\partial/\partial y_2)(\partial/\partial x_3)$, we write $\beta\gamma\delta$ and $x y x$, respectively.

GHS. The expression $\beta\gamma\delta$ goes over to

$$\frac{1}{8}(-x + y - z + w)(-x - y + z + w)(x - y - z + w) \equiv T_1 + T_2 + T_3, \quad (A.1)$$

where the terms in T_1 contain only a single variable, those in T_2 two different variables, and those in T_3 three different variables. We see that

$$T_1 = \frac{1}{8}(x x x + y y y + z z z + w w w), \quad (A.2)$$

which goes over to $xxx/2$ after permutation. The terms in T_2 fall into three separate groups: those where the repeated variable appears in the first and second slot, in the first and third slot, and in the second and third slot, respectively. Taking into account the signs and permuting, we see that the three groups contribute $-xxy/2$, $-xyx/2$, and $-xyy/2$, respectively. Finally, T_3 contributes xyz . Thus, $2\beta\gamma\delta$ goes over to

$$xxx - xxy - xyx - xyy + 2xyz, \quad (\text{A.3})$$

which is D in (3.5).

GKS. $\tilde{\mathcal{G}} = \partial/\partial\alpha_1$ goes over to $(x + y + z + w)/2$, which is equivalent to $2x$. Hence,

$$c_{\mathcal{G}} = 2\langle\sigma_1\rangle. \quad (\text{A.4})$$

$\tilde{\mathcal{H}} = (\partial/\partial\beta_1)(\partial/\partial\beta_2)$ goes over to $(-x + y - z + w)(-x + y - z + w)/4$, which is equivalent to $xx - xy$. Hence

$$c_{\mathcal{H}} = \langle\sigma_1\sigma_2\rangle - \langle\sigma_1\rangle\langle\sigma_2\rangle. \quad (\text{A.5})$$

Since we are in case b) of Theorem 1, the expressions in (A.4) and (A.5) are both non-negative.

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Note Added in Proof. One consequence of the GHS inequality is the negativity at $h=0$ of the Fourth Ursell Function $u_4(i, j, k, l)$. This also follows directly from Theorem 1. Indeed, $\tilde{\mathcal{H}} = (\partial/\partial\alpha_i)(\partial/\partial\beta_j)(\partial/\partial\gamma_k)(\partial/\partial\delta_l)$ can be shown to give rise to $u_4(i, j, k, l)$ at $h=0$. Since this $\tilde{\mathcal{H}}$ is covered by case a) of Theorem 1, the negativity of u_4 follows.