

HW #1 Solutions

1. *Fourier Series and Orthogonality of Sines and Cosines:* The Fourier series for  $f \in L^2[-\pi, \pi]$ , given by

$$f(x) = \frac{a_0}{2} + \sum_{k=1}^{\infty} a_k \cos kx + b_k \sin kx,$$

is an expansion of the function  $f(x)$  in the basis of trigonometric functions

$$\{\phi_0(x), \phi_1(x), \phi_2(x), \phi_3(x), \phi_4(x), \dots\} = \{1, \cos x, \sin x, \cos 2x, \sin 2x, \dots\}$$

Recall the inner product

$$\langle f, g \rangle = \int_{-\pi}^{\pi} f(x)g(x) dx,$$

find an expression for  $\langle \phi_n(x), \phi_m(x) \rangle$  for  $m, n \geq 0$ . In particular, show that the basis functions are mutually orthogonal by showing that for  $n \neq m$  the inner product is zero.

**ANS:**

Using the identity  $\sin a \sin b = \frac{1}{2}(\cos(a - b) - \cos(a + b))$

$$\begin{aligned} \langle \sin nx, \sin mx \rangle &= \int_{-\pi}^{\pi} \sin nx \sin mx dx \\ &= \frac{1}{2} \int_{-\pi}^{\pi} \cos(n - m)x - \cos(n + m)x dx \\ &= \frac{1}{2} \begin{cases} \int_{-\pi}^{\pi} 1 - 1 dx & n = m = 0 \\ \int_{-\pi}^{\pi} 1 - \cos 2nx dx & n = m \neq 0 \\ \int_{-\pi}^{\pi} \cos(n - m)x - \cos(n + m)x dx & n \neq m \end{cases} \\ &= \frac{1}{2} \begin{cases} 0 & n = m = 0 \\ 2\pi - \left[ \frac{1}{2n} \sin 2nx \right]_{-\pi}^{\pi} & n = m \neq 0 \\ \left[ \frac{1}{(n-m)} \sin(n - m)x - \frac{1}{(n+m)} \sin(n + m)x \right]_{-\pi}^{\pi} & n \neq m \end{cases} \\ &= \frac{1}{2} \begin{cases} 0 & n = m = 0 \\ 2\pi - 0 & n = m \neq 0 \\ 0 - 0 & n \neq m \end{cases} \\ &= \begin{cases} 0 & n = m = 0 \\ \pi & n = m \neq 0 \\ 0 & n \neq m \end{cases} \end{aligned}$$

Using the identity  $\cos a \cos b = \frac{1}{2}(\cos(a - b) + \cos(a + b))$  and following steps similar to those above:

$$\begin{aligned} \langle \cos nx, \cos mx \rangle &= \int_{-\pi}^{\pi} \cos nx \cos mx \, dx \\ &= \frac{1}{2} \int_{-\pi}^{\pi} \cos(n - m)x + \cos(n + m)x \, dx \\ &= \begin{cases} 2\pi & n = m = 0 \\ \pi & n = m \neq 0 \\ 0 & n \neq m \end{cases} \end{aligned}$$

Using the identity  $\sin a \cos b = \frac{1}{2}(\sin(a - b) + \sin(a + b))$ . Note below that if  $n = m = 0$  or  $m = 0$  and  $n > 0$  the inner product is easily seen to be 0. Thus, assuming  $m, n > 0$

$$\begin{aligned} \langle \sin nx, \cos mx \rangle &= \int_{-\pi}^{\pi} \sin nx \cos mx \, dx \\ &= \frac{1}{2} \int_{-\pi}^{\pi} \sin(n - m)x + \sin(n + m)x \, dx \\ &= \begin{cases} \int_{-\pi}^{\pi} 0 + \sin 2nx \, dx & n = m \\ - \left[ \frac{1}{(n-m)} \cos(n - m)x + \frac{1}{(n+m)} \cos(n + m)x \right] \Big|_{-\pi}^{\pi} & n \neq m \end{cases} \\ &= 0. \end{aligned}$$

2. *Separation of Variables*: Use separation of variables to find the solution, in the form of an infinite series, of the homogeneous heat conduction problem with Neumann no flux boundary conditions:

$$\begin{aligned} \text{PDE: } & \frac{\partial u}{\partial t} = \kappa \frac{\partial^2 u}{\partial x^2}, \quad (0 < x < L, t > 0) \\ \text{BCs: } & \frac{\partial u}{\partial x}(0, t) = 0, \quad \frac{\partial u}{\partial x}(L, t) = 0, \quad (t > 0) \\ \text{IC: } & u(x, 0) = f(x) \quad (t = 0) \end{aligned}$$

Proceed as follows:

- Assume  $u(x, t) = G(t)\phi(x)$  and derive the ODEs satisfied by  $\phi(x)$  and  $G(t)$ .
- Solve the ODEs for  $\phi(x)$  and  $G(t)$ , and determine the allowed values for the separation constant  $\lambda$ .
- Show that the eigenfunctions of the spatial eigenvalue-eigenfunction problem are mutually orthogonal.
- Write the solution in terms of an infinite series with coefficients  $a_n$ , and derive a formula for the  $a_n$  in terms of an integral involving the initial condition  $u(x, 0) = f(x)$ .

**ANS:** To solve using separation of variables, we proceed as follows:

- Letting  $u(x, t) = G(t)\phi(x)$  and substituting this expression into the PDE gives

$$G'(t)\phi(x) = \kappa G(t)\phi''(x) \quad \text{or} \quad \frac{G'(t)}{\kappa G(t)} = \frac{\phi''(x)}{\phi(x)} = -\lambda$$

where  $\lambda$  is the separation constant. Therefore  $G(t)$  satisfies the ODE

$$G'(t) = -\lambda\kappa G(t),$$

and  $\phi(x)$  the ODE with BC

$$\begin{cases} \phi''(x) = -\lambda\phi(x) \\ \phi'(0) = \phi'(L) = 0. \end{cases}$$

- The ODE for  $G(t)$  has the general solution

$$G(t) = Ce^{-\kappa\lambda t},$$

while the solution  $\phi(x)$  depends on  $\lambda$ :

$\lambda < 0$ : General solution  $\phi(x) = c_1 \cosh \sqrt{-\lambda}x + c_2 \sinh \sqrt{-\lambda}x$ , and applying BC

$$0 = \phi'(0) = \sqrt{-\lambda} (c_1 \sinh 0 + c_2 \cosh 0) = \sqrt{-\lambda} (c_1 * 0 + c_2 * 1) = \sqrt{-\lambda} c_2 \quad \Rightarrow \quad c_2 = 0,$$

and

$$0 = \phi'(L) = c_1 \sqrt{-\lambda} \sinh \sqrt{-\lambda} L \Rightarrow c_1 = 0 \text{ since } \sinh \sqrt{-\lambda} L \neq 0.$$

Thus for  $\lambda < 0$  we have only the trivial solution  $\phi \equiv 0$ .

$\lambda = 0$ : General solution  $\phi(x) = c_1 + c_2 x$ , and applying BC

$$0 = \phi'(0) = \phi'(L) = c_2 \Rightarrow c_2 = 0.$$

Thus for  $\lambda = 0$  we have  $\phi = c_1 = A_0$  where  $A_0$  is an arbitrary nonzero number, which we can set to 1.

$\lambda > 0$ : General solution  $\phi(x) = c_1 \cos \sqrt{\lambda} x + c_2 \sin \sqrt{\lambda} x$ , and applying BC

$$0 = \phi'(0) = \sqrt{\lambda} (-c_1 \sin 0 + c_2 \cos 0) = \sqrt{\lambda} (-c_1 * 0 + c_2 * 1) = \sqrt{\lambda} c_2 \Rightarrow c_2 = 0,$$

and

$$0 = \phi'(L) = -\sqrt{\lambda} c_1 \sin \sqrt{\lambda} L \Rightarrow c_1 = 0 \text{ or } \sqrt{\lambda} L = n\pi \text{ for } n = 1, 2, 3, \dots$$

Since  $c_1 = 0$  gives the trivial solution, we take the latter, and have for  $\lambda > 0$

$$\lambda_n = \frac{n\pi^2}{L^2}, \quad \phi_n(x) = A_n \cos \frac{n\pi x}{L} \text{ for each } n = 1, 2, 3, \dots$$

(c) The eigenfunctions of the spatial eigenvalue-eigenfunction problem are given by

$$\{\phi_0, \phi_1, \phi_2, \dots\} = \left\{1, \cos \frac{\pi x}{L}, \cos \frac{2\pi x}{L}, \dots\right\} = \{\phi_n(x)\}_{n=0}^{\infty}.$$

Using the results of problem 2, we have for  $n, m \geq 0$

$$\begin{aligned} \langle \phi_n(x), \phi_m(x) \rangle &= \int_0^L \cos \frac{n\pi x}{L} \cos \frac{m\pi x}{L} dx \\ &= \frac{1}{2} \int_0^L \cos \frac{(n-m)\pi x}{L} + \cos \frac{(n+m)\pi x}{L} dx \\ &= \begin{cases} L & n = m = 0 \\ L/2 & n = m \neq 0 \\ 0 & n \neq m \end{cases} \end{aligned}$$

showing that when  $n \neq m$  that  $\langle \phi_n(x), \phi_m(x) \rangle = 0$ , i.e. the set of eigenfunctions  $\{\phi_n(x)\}_{n=0}^{\infty}$  is mutually orthogonal.

(d) Using the orthogonality of the  $\phi_n$ , we have

$$u(x, t) = A_0 + \sum_{n=1}^{\infty} A_n e^{-\kappa \left(\frac{n\pi}{L}\right)^2 t} \cos \frac{n\pi x}{L}$$

where

$$A_0 = \frac{1}{L} \int_0^L f(x) dx \quad \text{and} \quad A_n = \frac{2}{L} \int_0^L f(x) \cos \frac{n\pi x}{L} dx.$$

3. *Inhomogeneous Boundary Conditions:* Consider the 1D heat conduction problem for  $u(x, t)$  with fixed temperature boundary conditions:

$$\begin{aligned} \text{PDE: } & \frac{\partial u}{\partial t} = \kappa \frac{\partial^2 u}{\partial x^2}, \quad (0 < x < L, t > 0) \\ \text{BCs: } & u(0, t) = \alpha, \quad u(L, t) = \beta, \quad (t > 0) \\ \text{IC: } & u(x, 0) = f(x) \quad (t = 0) \end{aligned}$$

Let  $h(x) = \alpha + (\beta - \alpha)x/L$  and  $v(x, t) = u(x, t) - h(x)$ . Determine the PDE, BCs and IC that  $v(x, t)$  satisfies. Given the solution  $v(x, t)$  of the new system, explain the steps that you would go about to solve the original inhomogeneous problem for  $u(x, t)$ .

ANS: Let  $v(x, t) = u(x, t) - h(x)$ . Then

$$\frac{\partial v}{\partial t} = \frac{\partial u}{\partial t} \quad \text{and} \quad \frac{\partial^2 v}{\partial x^2} = \frac{\partial^2 u}{\partial x^2} \quad \text{since} \quad \frac{d}{dt}h(x) = \frac{d^2}{dx^2}h(x) = 0,$$

and note that  $h(0) = \alpha$  and  $h(L) = \beta$  giving  $v(0, t) = v(L, t) = 0$ . Then the system for  $v(x, t)$  is given by

$$\begin{cases} \frac{\partial v}{\partial t} = \kappa \frac{\partial^2 v}{\partial x^2} & (0 < x < L, t > 0) \\ v(0, t) = v(L, t) = 0 & (t > 0) \\ v(x, 0) = u(x, 0) - h(x) = f(x) - h(x) = g(x) & (t = 0) \end{cases}$$

Thus to solve the original system for  $u(x, t)$ , first solve the system for  $v(x, t)$  and one has  $u(x, t) = v(x, t) + h(x)$ .

4. *Self-Adjointness of the Laplacian:* Consider the vector space  $V = C^2([a, b])$ , the set of all real-valued functions  $f(x)$  defined on the interval  $[a, b]$  which are at least two times continuously differentiable. Let the inner product on  $V$  be defined in the usual manner,  $\langle f, g \rangle = \int_a^b f(x)g(x) dx$ . Show that for any  $f, g \in V$  that satisfy the boundary conditions

- (a)  $f(a) = f(b) = 0$ , and similarly for  $g$  (Dirichlet BC).
- (b)  $f'(a) = f'(b) = 0$ , and similarly for  $g$  (Neumann BC).
- (c)  $f(a) = f(b)$ ,  $f'(a) = f'(b)$ , and similarly for  $g$  (periodic BC).

that, considering each case above **separately**,

$$\left\langle \frac{d^2}{dx^2}f, g \right\rangle = \left\langle f, \frac{d^2}{dx^2}g \right\rangle,$$

This shows that the 1D Laplacian, considered as a linear operator on  $V$  along with each of the above boundary conditions, is a *self-adjoint* operator. (Hint: In each case, use integration by parts for definite integrals *twice* and apply the boundary conditions.)

**ANS:** Now, integrating by parts twice, and using the fact that the boundary terms that arise in each of the two steps is 0 for each of the BCs listed above, we have

$$\begin{aligned} \left\langle \frac{d^2}{dx^2}f, g \right\rangle &= \int_a^b f''(x)g(x) dx \quad \text{and integrating by parts} \\ &= g(x)f'(x) \Big|_a^b - \int_a^b f'(x)g'(x) dx \\ &= - \int_a^b f'(x)g'(x) dx \quad \text{and integrating by parts} \\ &= - \left[ g'(x)f(x) \Big|_a^b - \int_a^b f(x)g''(x) dx \right] \\ &= \int_a^b f(x)g''(x) dx \\ &= \left\langle f, \frac{d^2}{dx^2}g \right\rangle \end{aligned}$$