Metaplectic Whittaker Functions and Crystals of Type B

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September 8, 2010

Let n be an integer and let F be a nonarchimedean local field whose characteristic is not a prime dividing n. Let μ_k be the group of k-th roots of unity in the algebraic closure of F; we assume that $\mu_{2n} \subset F$. Let G be a split, simply-connected semisimple algebraic group over F. We assume that G is actually defined over the ring $\mathfrak{o}$ of integers in F in such a way that $K = G(\mathfrak{o})$ is a special maximal compact subgroup of $G(F)$.

Matsumoto [20] constructed an n-fold metaplectic cover $\tilde{G}(F)$ of $G(F)$. For this, we only need $\mu_n \subset F$ but the hypothesis $\mu_{2n} \subset F$ simplifies the metaplectic cocycle and the resulting formulas. We are interested in values of a spherical Whittaker function W on $\tilde{G}(F)$.

Let $G = \mathrm{Sp}_{2r}$ and let the cover degree $n = 2$. In this case, we connect a known description of the Whittaker function to the theory of multiple Dirichlet series:

- Bump, Friedberg and Hoffstein [8] gave a description of the Whittaker function, essentially as a sum of at most 2^r irreducible characters of $\mathrm{Sp}(2r)$, that is, of Cartan type C_r.

- Chinta and Gunnells [14] gave a recipe for the p-parts of Weyl group multiple Dirichlet series for any root system Φ and any positive integer n. In the special case $\Phi = B_r$ and $n = 2$, we show that this agrees with the description of the Whittaker function in [8].

In addition to these descriptions, we have three other conjectural formulas for the metaplectic Whittaker function on Sp_{2r} using the root system of type B_r. Let λ denote a dominant weight for this root system and let t_λ be an element of the split maximal torus parametrized by λ. Then:

- The value of the Whittaker function W at t_λ may be expressed as a sum over the Kashiwara crystal $\mathcal{B}_\lambda$. (Conjecture 1)
- The value $W(t_\lambda)$ may be expressed as a sum over the Kashiwara crystal $\mathcal{B}_{\lambda+\rho}$, where ρ is the Weyl vector. (Conjecture 2)
- The value $W(t_\lambda)$ may be expressed as the partition function for a statistical lattice model – square ice with U-turn boundary. (Conjecture 3)

The second and third conjectural descriptions are easily seen to be equivalent but give rise to very different considerations. We offer partial proofs of these conjectures in the following sections. The conjectures are further convincingly supported by extensive calculations using SAGE.

An interesting feature of this situation is the interplay between Type B descriptions and Type C descriptions.

The Classical Case: The Casselman-Shalika formula

Before considering the metaplectic case, let us review the situation when $n = 1$, so that $G(F)$ and $\tilde{G}(F)$ are the same. Let Λ be the weight lattice of the connected L-group ${}^LG^\circ$. It is the group $X({}^LT)$ of rational characters of a maximal torus ${}^LT$ of ${}^LG^\circ$. If $\lambda \in \Lambda$ and $\boldsymbol{z} \in {}^LT$ we will denote by $\boldsymbol{z}^\lambda$ the value of λ at $\boldsymbol{z}$. Let Φ be the root system of ${}^LG^\circ$, so that the root system of G is the dual root system $\hat{\Phi}$.

If T is an F-split torus of G, then $\Lambda \cong T(F)/T(\mathfrak{o})$. If $\lambda \in \Lambda$, let t_λ be a representative of its coset in $T(F)$. Unramified quasicharacters of $T(F)$ correspond to elements of ${}^LT$. Indeed, an unramified quasicharacter ξ of $T(F)$ is a quasicharacter that is trivial on $T(\mathfrak{o})$, that is, a character of Λ, and so there is an element $\boldsymbol{z} \in {}^LT$ such that $\xi(t_\lambda) = \boldsymbol{z}^\lambda$. In this case, we write $\xi = \xi_{\boldsymbol{z}}$.

If α is a positive root, then the coroot $\alpha^\vee$ is a positive root of G with respect to T. Let $X_{\alpha^\vee}$ be the corresponding root eigenspace in Lie(G), and let N be the maximal unipotent subgroup with Lie algebra $\bigoplus_{\alpha\in\Phi^+} X_{\alpha^\vee}$. Then $B = TN$ is a Borel subgroup.

Let ψ_N be a nondegenerate character of N. Then ψ_N is trivial on $\exp(X_{\alpha^\vee})$ if α is positive root that is not simple. If α is a simple positive root then we may arrange that ψ_N is trivial on $\exp(X_{\alpha^\vee}) \cap K$ but no larger subgroup of $\exp(X_{\alpha^\vee})$.

Let $\xi = \xi_{\boldsymbol{z}}$ be a character of $T(F)$, which we extend to a character of $B(F)$ by taking $N(F)$ to be in the kernel. Let δ be the modular quasicharacter of $B(F)$. The normalized induced representation $\pi(\xi)$ consists of all locally constant functions $f : G(F) \longrightarrow \mathbb{C}$ such that $f(bg) = (\xi\delta^{1/2})(b)f(g)$, with $G(F)$ acting by right translation.

The standard spherical vector f° is the unique function such that $f^\circ(k) = 1$ for $k \in K$. Let w_0 be a representative of the long Weyl group element. We may assume that $w_0 \in K$. Then the spherical Whittaker function is

$$W(g) = \int_{N(F)} f^\circ(w_0 ng)\psi_N(n)\, dn. \tag{1}$$

If $\xi = \xi_{\boldsymbol{z}}$ then the integral is convergent provided $|\boldsymbol{z}^\alpha| < 1$ for $\alpha \in \Phi^+$. For other $\boldsymbol{z}$ it may be defined by analytic continuation from this domain.

According to the formula of Casselman and Shalika [9] we have $W(t_\lambda) = 0$ unless the weight λ is dominant, and if λ is dominant, then

$$W(t_\lambda) = \prod_{\alpha\in\Phi^+} (1 - q^{-1}\boldsymbol{z}^\alpha)\chi_\lambda(\boldsymbol{z}), \tag{2}$$

where χ_λ is the irreducible character of ${}^L G^\circ$ with highest weight λ and q is the cardinality of the residue field.

Let $\mathcal{B}_\lambda$ be the Kashiwara crystal with highest weight λ, so that

$$\chi_\lambda(\boldsymbol{z}) = \sum_{v\in\mathcal{B}_\lambda} \boldsymbol{z}^{\mathrm{wt}(v)}. \tag{3}$$

Ignoring the normalizing constant $\prod_{\alpha\in\Phi^+}(1 - q^{-1}\boldsymbol{z}^\alpha)$ in (2), this could be regarded as a formula for the Whittaker function.

We note that by the Weyl character formula

$$\prod_{\alpha\in\Phi^+} (1 - \boldsymbol{z}^\alpha)\chi_\lambda(\boldsymbol{z}) = \sum_{w\in W} (-1)^{l(w)} \boldsymbol{z}^{w(\rho+\lambda)+\rho}, \qquad \rho = \frac{1}{2}\sum_{\alpha\in\Phi^+} \alpha.$$

The factor $\prod_{\alpha\in\Phi^+}(1-\boldsymbol{z}^\alpha)$ is the *Weyl denominator* and the factor $\prod_{\alpha\in\Phi^+}(1-q^{-1}\boldsymbol{z}^\alpha)$ which appears in (2) is a deformation of this factor.

We are therefore interested in deformations of the Weyl character formula in which the deformed denominator appears. A typical such formula will have the form

$$\prod_{\alpha\in\Phi^+} (1 - q^{-1}\boldsymbol{z}^\alpha)\chi_\lambda(\boldsymbol{z}) = \sum_{v\in\mathcal{B}_{\lambda+\rho}} G(v)\boldsymbol{z}^{\mathrm{wt}(v)}, \tag{4}$$

where $\mathcal{B}_{\lambda+\rho}$ is the Kashiwara crystal with highest weight $\lambda+\rho$. We will call a function G on $\mathcal{B}_{\lambda+\rho}$ which satisfies this identity a *Tokuyama function*. The archetype is the formula of Tokuyama [22], where it was stated in the language of Gelfand-Tsetlin patterns, and translated into the crystal language in [5]. This is for Cartan type A. For Cartan types C and D, see [1], [10] in this volume.

For general n, we may define the metaplectic Whittaker function by an integral generalizing (1), and then ask for a formula of the form

$$W(t_\lambda) = \delta^{1/2}(t_\lambda) \sum_{v \in \mathcal{B}_{\lambda+\rho}} G(v) \boldsymbol{z}^{\mathrm{wt}(v)}. \tag{5}$$

We will give analogs of both (3) and (4) for the metaplectic Whittaker function on the double cover of $\mathrm{Sp}_{2r}(F)$. However this is the *only* metaplectic example where we have an analog of (3), whereas analogs of (4) may be found in many cases of group and degree of metaplectic cover:

- $G = \mathrm{SL}_n$ and any n: [5], [6], [7].
- $G = \mathrm{Spin}(2r+1)$ and n odd: [1] (rigorously for $n = 1$ or n sufficiently large).
- $G = \mathrm{Spin}(2r)$ and n even: [10].
- $G = \mathrm{Sp}(2r)$ and n even: this paper (rigorously for $n = 2$).

The Metaplectic Whittaker Function

We review the formula for the metaplectic Whittaker function on the double cover of $\mathrm{Sp}_{2r}(F)$ which was found by Bump, Friedberg and Hoffstein. We are assuming that $\mu_4 \subset F$, which simplifies the formula slightly, since the quadratic Hilbert symbol $(-1, a)_2 = (a, a)_2 = 1$ because -1 is a square.

Let $\mathrm{Sp}_{2r} = \{g \in \mathrm{GL}_{2r} \mid gJ^tg = J\}$, where $J = \begin{pmatrix} & -J_r \\ J_r & \end{pmatrix}$, $J_r = \begin{pmatrix} & & 1 \\ & \iddots & \\ 1 & & \end{pmatrix}$.

The metaplectic cocycle defining the double cover satisfies

$$\sigma\left(\begin{pmatrix} x_1 & & & & & \\ & \ddots & & & & \\ & & x_r & & & \\ & & & x_r^{-1} & & \\ & & & & \ddots & \\ & & & & & x_1^{-1} \end{pmatrix}, \begin{pmatrix} y_1 & & & & & \\ & \ddots & & & & \\ & & y_r & & & \\ & & & y_r^{-1} & & \\ & & & & \ddots & \\ & & & & & y_1^{-1} \end{pmatrix}\right) = \prod (x_i, y_i)_2.$$

The double cover $\widetilde{\mathrm{Sp}}_{2r}(F)$ consists of pairs (g, ε) with $g \in \mathrm{Sp}_{2r}(F)$ and $\varepsilon = \pm 1$. The multiplication is $(g, \varepsilon)(g', \varepsilon') = (gg', \varepsilon\varepsilon'\sigma(g, g'))$. Let $\Lambda_C = \mathbb{Z}^r$; in the next section we will interpret this as the weight lattice of Cartan Type C_r. An element $\lambda = (\lambda_1, \cdots, \lambda_r) \in \Lambda_C$ is *dominant* if $\lambda_1 \geqslant \ldots \geqslant \lambda_r \geqslant 0$. We define the "alternator"

$$\mathcal{A} = \sum_{w \in W} (-1)^{l(w)} w \tag{6}$$

as a member of the group algebra of the Weyl group W. As a group acting on the spectral parameters $\boldsymbol{z} = (z_1, \cdots, z_r)$, W is the group generated by the $r!$ permutations, and the 2^r transformations $z_i \to z_i^{\pm 1}$. The r simple reflections $s_1, \ldots, s_r \in W$ correspond to $s_i \colon z_i \leftrightarrow z_{i+1}$ for $i = 1, \ldots, r-1$, and $s_r \colon z_r \mapsto 1/z_r$. We will denote $\boldsymbol{z}^\lambda = \prod z_i^{\lambda_i}$ for $\lambda \in \Lambda_C$. Let $\rho_C = (r, r-1, \cdots, 1)$ denote the Weyl vector. By the Weyl denominator formula

$$\Delta_C(\boldsymbol{z}) := \sum_{w \in W} (-1)^{\ell(w)} w(\boldsymbol{z}^{\rho_C}) = \boldsymbol{z}^{-\rho_C} \prod_{i=1}^{r} (1 - z_i^2) \prod_{i<j} (1 - z_i z_j)(1 - z_i z_j^{-1}).$$

We sometimes simply write the denominator as Δ_C, when clear from context.

If $\lambda \in \Lambda_C$, let

$$t_\lambda = \begin{pmatrix} p^{\lambda_1} & & & & & \\ & \ddots & & & & \\ & & p^{\lambda_r} & & & \\ & & & p^{-\lambda_r} & & \\ & & & & \ddots & \\ & & & & & p^{-\lambda_1} \end{pmatrix}.$$

We fix an additive character ψ on F. This gives rise to a nondegenerate character ψ_N on the subgroup $N(F)$ of upper triangular unipotent matrices n of $\mathrm{Sp}_{2r}(F)$ by $\psi_N(n) = \psi(n_{12} + n_{23} + \cdots n_{r,r+1})$. The cocycle $\sigma(n, g) = \sigma(g, n) = 1$ for $n \in N(F)$ and g arbitrary, so the map $N(F) \longrightarrow \widetilde{\mathrm{Sp}}_{2r}(F)$ given by $n \mapsto (n, 1)$ is a homomorphism, and we may identify $N(F)$ with its image.

If $a \in F^\times$, let $\gamma(a) = \sqrt{|a|} \int \psi(ax^2)\, dx / \int \psi(x^2)\, dx$ where the integral is taken over any sufficiently large fractional ideal. Let $\boldsymbol{s} : T(F) \longrightarrow \widetilde{\mathrm{Sp}}_{2r}(F)$ be the map $t \mapsto \boldsymbol{s}(t) = (t, 1)$. Then $\gamma(ab)/\gamma(a)\gamma(b) = (a, b)_2$, the local Hilbert symbol.

Theorem 1 (Bump, Friedberg, Hoffstein) *If $\lambda \in \Lambda_C$ is dominant, we have*

$$W(t_\lambda) = \delta^{1/2}(t_\lambda) \frac{1}{\Delta_C} \mathcal{A} \left(\boldsymbol{z}^{\lambda + \rho_C} \prod_{k=1}^{r} (1 - q^{-1/2} z_i^{-1}) \right) W(1).$$

Moreover

$$W(1) = \left(\prod_{i=1}^{r} \gamma(p^{\lambda_i})^{-1}\right) \prod_{i} (1 + q^{-\frac{1}{2}} z_i) \prod_{i<j} (1 - q^{-1} z_i z_j^{-1})(1 - q^{-1} z_i z_j).$$

If λ is not dominant then $W(t_\lambda) = 0$.

Let us combine the two most important parts of this formula and write

$$\begin{aligned} W(\lambda) \;&=\; \prod_{i} (1 + q^{-\frac{1}{2}} z_i) \prod_{i<j} (1 - q^{-1} z_i z_j^{-1})(1 - q^{-1} z_i z_j) \times \\ &\quad \frac{1}{\Delta_C} \mathcal{A}\left(\boldsymbol{z}^{\lambda+\rho_C} \prod_{k=1}^{r} (1 - q^{-1/2} \boldsymbol{z}_i^{-1})\right). \end{aligned} \tag{7}$$

We note that in this context λ is integral but (7) makes sense if λ is half-integral. Furthermore, the Whittaker function can be extended to the larger group GSp_{2r}. It is natural to expect that our results can be extended to GSp_{2r}, and that the values of (7) when λ is half-integral are to be interpreted as values of the Whittaker function on GSp_{2r}. Although we cannot confirm this when λ is half-integral, we will make some observations about the values of (7) in this case.

An Embarrassment of L-groups

Although Langlands only defined an L-group for algebraic groups, there is a natural candidate for an L-group of $\tilde{G}(F)$ when G is split. For $G = \mathrm{Sp}_{2r}$, it is natural to assume that the L-group should be:

$$\begin{cases} \mathrm{Sp}_{2r}(\mathbb{C}) & \text{if } n \text{ is even,} \\ \mathrm{Spin}_{2r+1}(\mathbb{C}) & \text{if } n \text{ is odd.} \end{cases}$$

For example, the alternation of the Cartan type of the L-group is suggested by Savin [21], who found that the Cartan type of the genuine part of the Iwahori Hecke algebra was isomorphic to that of $\mathrm{Sp}_{2r}(F)$ if n is odd and of $\mathrm{Spin}_{2r+1}(F)$ if n is even, suggesting that the L-group of the metaplectic should be isomorphic to the L-groups of these groups. Thus we may provisionally expect that in generalizing the Casselman-Shalika formula to the double cover of Sp_{2r} the role of ${}^LG^\circ$ should be played by $\mathrm{Sp}_{2r}(\mathbb{C})$, and indeed, such a generalization was found by Bump, Friedberg and Hoffstein [8].

It is therefore a little surprising that in generalizing (5) the relevant crystal $\mathcal{B}_\lambda$ is not of type C_r but rather of type B_r! In explaining this, both the representations of $\mathrm{Sp}_{2r}(\mathbb{C})$ (type C_r) and $\mathrm{Spin}_{2r+1}(\mathbb{C})$ (type B_r) will play a role.

We will compare these representation theories by the *ad hoc* method of identifying the ambient spaces of their weight lattices. The weight lattice Λ_C of type C_r is $\mathbb{Z}^r$. The lattice Λ_C has index two in the weight lattice Λ_B of type B_r. The lattice Λ_B consists of $\lambda = (\lambda_1, \cdots, \lambda_r) \in \frac{1}{2}\mathbb{Z}^r$ such that all $\lambda_i - \lambda_j \in \mathbb{Z}$. The Weyl group W of type B_r is the same as the Weyl group of type C_r; acting on Λ_B or Λ_C, it is generated by simple reflections $s_1, \cdots, s_r$ where s_i acting on $\Lambda = \mathbb{Z}^r$ interchanges λ_i and λ_{i+1} in $\lambda = (\lambda_1, \cdots, \lambda_r)$ when $i < r$, and s_r sends $\lambda_r \to -\lambda_r$. The Weyl vector ρ of any root system is half the sum of the positive roots and so for B_r and C_r, respectively, we have

$$\rho_B = \left(r - \frac{1}{2}, r - \frac{3}{2}, \cdots, \frac{1}{2} \right), \qquad \rho_C = (r, r-1, \cdots, 1).$$

If $\lambda \in \Lambda_C$ is a dominant weight, then the irreducible character of $\mathrm{Sp}_{2r}(\mathbb{C})$ with highest weight λ will be denoted χ_λ^C, and similarly if $\lambda \in \Lambda_B$ is a dominant weight, the irreducible character of $\mathrm{Spin}_{2r+1}(\mathbb{C})$ with highest weight λ will be denoted χ_λ^B. In either case, let g be an element of the relevant group. Let $\boldsymbol{z} = (z_1, \cdots, z_r)$ be such that the eigenvalues of g are $z_i^{\pm 1}$ in the symplectic case, or such that the eigenvalues of the image of g in $\mathrm{SO}_{2r+1}(\mathbb{C})$ are $z_i^{\pm 1}$ and 1 in the spin case. Then the Weyl character formula asserts that

$$\chi_\lambda^C(g) = \frac{1}{\Delta_C} \mathcal{A}(\boldsymbol{z}^{\rho_C + \lambda}) \qquad \text{or} \qquad \chi_\lambda^B(g) = \frac{1}{\Delta_B} \mathcal{A}(\boldsymbol{z}^{\rho_B + \lambda})$$

depending on which case we are in, where the Weyl denominators are

$$\Delta_C = \mathcal{A}(\boldsymbol{z}^{\rho_C}) = \prod_{i<j} \left[(z_i^{1/2} z_j^{-1/2} - z_i^{-1/2} z_j^{1/2})(z_i^{1/2} z_j^{1/2} - z_i^{-1/2} z_j^{-1/2}) \right] \prod_i (z_i - z_i^{-1}),$$

$$\Delta_B = \mathcal{A}(\boldsymbol{z}^{\rho_B}) = \prod_{i<j} \left[(z_i^{1/2} z_j^{-1/2} - z_i^{-1/2} z_j^{1/2})(z_i^{1/2} z_j^{1/2} - z_i^{-1/2} z_j^{-1/2}) \right] \prod_i (z_i^{1/2} - z_i^{-1/2}).$$

In particular

$$\frac{\Delta_C}{\Delta_B} = \prod_{i=1}^r (z_i^{1/2} + z_i^{-1/2}) = \frac{\rho_C}{\rho_B} \prod_{i=1}^r (1 - z_i^{-1}). \tag{8}$$

On the face of it, the last formula has little meaning, since the Weyl denominators live on different groups. We will use it in the next section.

Ambivalence of the L-group

Let G be a reductive group over a nonarchimedean local field F. Let us consider the role of the L-group in the Casselman-Shalika formula. The semisimple conjugacy classes of ${}^LG^\circ$ parametrize the spherical representations of $G(F)$. Let π be a spherical representation and $\boldsymbol{z} = \boldsymbol{z}_\pi$ the parametrizing conjugacy class. Then the values of the irreducible characters of $G(F)$ on $\boldsymbol{z}$ equal the values of the spherical Whittaker function of π.

So we should seek a similar interpretation in the metaplectic case. Let $G = \mathrm{Sp}_{2r}(F)$ and let $\tilde{G}(F)$ be the double cover. Either $\mathrm{Sp}_{2r}(\mathbb{C})$ or $\mathrm{SO}_{2r+1}(\mathbb{C})$ will serve to parametrize the principal series representations of G.

We first seek an interpretation of the factor

$$\prod_i (1 + q^{-\frac{1}{2}} z_i) \prod_{i<j} (1 - q^{-1} z_i z_j^{-1})(1 - q^{-1} z_i z_j) \tag{9}$$

appearing in (7) as a deformation of a Weyl denominator. The Weyl denominators of types B and C are, respectively, $\boldsymbol{z}^{-\rho_B}$ and $\boldsymbol{z}^{-\rho_C}$ times

$$\prod_i (1 - z_i) \prod_{i<j} (1 - z_i z_j^{-1})(1 - z_i z_j), \qquad \text{and} \qquad \prod_i (1 - z_i^2) \prod_{i<j} (1 - z_i z_j^{-1})(1 - z_i z_j).$$

Now there are two ways of looking at (9). We may write it as

$$\prod_i (1 - q^{-\frac{1}{2}} z_i)^{-1} \times \prod_i (1 - q^{-1} z_i^2) \prod_{i<j} (1 - q^{-1} z_i z_j^{-1})(1 - q^{-1} z_i z_j),$$

and the factor in front may be interpreted as the p-part of a quadratic L-function. The remaining terms in the product give the typical deformation of the Weyl denominator of type C, and taking the classical limit $q^{-1} \to 1$ recovers the familiar denominator formula in type C. On the other hand, we may let $q^{-\frac{1}{2}} \to -1$, in which case (9) becomes the Weyl denominator of type B.

A similar dual interpretation pertains with the factor

$$\frac{1}{\Delta_C} \mathcal{A} \left(\boldsymbol{z}^{\lambda + \rho_C} \prod_{k=1}^{r} (1 - q^{-1/2} z_i^{-1}) \right). \tag{10}$$

On the one hand, if we expand the product we get a sum

$$\sum_{S \subset \{1,2,3,\cdots,r\}} (-q^{1/2})^{|S|} \frac{1}{\Delta_C} \mathcal{A} \left(\boldsymbol{z}^{\lambda + \rho_C} \prod_{i \in S} z_i^{-1} \right). \tag{11}$$

Each term is either zero, or an irreducible character of $\mathrm{Sp}_{2r}(\mathbb{C})$ by the Weyl character formula. Hence (10) may be regarded as a sum of $\leqslant 2^r$ irreducible characters of $\mathrm{Sp}_{2r}(\mathbb{C})$ and thus has a Type C flavor. But on the other hand, let us again specialize $q^{\frac{1}{2}} \to -1$. Then using (8), the factor (10) becomes

$$\frac{1}{\Delta_B} \mathcal{A}\left(\boldsymbol{z}^{\lambda+\rho_B}\right) = \chi_\lambda^B(\boldsymbol{z}). \tag{12}$$

This formula generalizes to a formula like (3) for the metaplectic Whittaker function in the form (10). We will discuss this point in a subsequent section.

The averaging method of Chinta-Gunnells

Chinta and Gunnells [13], [14] gave a construction of the p-parts for multiple Dirichlet series which applies to any root system and choice of fixed positive integer n. In this section, we show that their construction gives the metaplectic Whittaker function of $\widetilde{\mathrm{Sp}}_{2r}(F)$ when the root system is of type B_r and the cover degree $n = 2$. (For additional articles on relations between multiple Dirichlet series and Whittaker functions, see also Chinta and Offen [11] and Chinta, Friedberg and Gunnells [12].)

The Chinta-Gunnells method begins by defining an action on rational functions of the spectral parameters, which we review in the case at hand.

As before, let T be the maximal torus of diagonal elements in SO_{2r+1}, whose eigenvalues are $z_1, \cdots, z_r, 1, z_r^{-1}, \cdots, z_1^{-1}$. Let T' be the preimage of T in Spin_{2r+1}. The coordinate ring $\mathbb{C}[T']$ of T' is then generated by $z_i^{\pm 1}$ and by $\sqrt{z_1 \cdots z_r}$. We remark that $\mathbb{C}[T']$ can be identified with the group ring $\mathbb{C}[\Lambda_B]$, with the z_i, $1 \leq i < r$ corresponding to the first $r-1$ fundamental weights and the product $\sqrt{z_1 \cdots z_r}$ corresponding to the spin representation (as before we think of Λ_C sitting as a sublattice of Λ_B of index 2).

Let $\mathbb{C}(T')$ be the fraction field of $\mathbb{C}[T']$. Consider the rational map $\mathbb{C}(T') \to \mathbb{C}(T')$ that takes $z_i \to z_i$ for $i = 1, \ldots, r-1$ and $\sqrt{z_1 \cdots z_r} \to -\sqrt{z_1 \cdots z_r}$. We write this map as $f(\boldsymbol{z}) \mapsto f(\varepsilon \boldsymbol{z})$, slightly abusing notation from before. Then we define an action of W by

$$(f|s_i)(\boldsymbol{z}) = f(s_i \boldsymbol{z}), \quad 1 \leq i < r,$$

and

$$(f|s_r)(\boldsymbol{z}) = \frac{1 - q^{-1/2} z_r^{-1}}{1 - q^{-1/2} z_r} f^+(s_r \boldsymbol{z}) + \frac{1}{z_r} f^-(s_r \boldsymbol{z}),$$

where

$$f^+(\boldsymbol{z}) = \frac{f(\boldsymbol{z}) + f(\varepsilon \boldsymbol{z})}{2}, \quad f^-(\boldsymbol{z}) = \frac{f(\boldsymbol{z}) - f(\varepsilon \boldsymbol{z})}{2}.$$

The braid relations are satisfied, and so this definition extend to a right action $f \mapsto f|w$ for all $w \in W$. Now the Chinta-Gunnells description of the p-part of the multiple Dirichlet series may be written

$$\Xi(\lambda, \boldsymbol{z}) = \boldsymbol{z}^{\lambda+\rho_C} \sum_{w \in W} \frac{\boldsymbol{z}^{-\lambda-\rho_C}|w}{\Delta_C(w\boldsymbol{z})}. \tag{13}$$

We further define a q-deformation of the Weyl denominator Δ_C by

$$D(\boldsymbol{z}; q) = \prod_{i=1}^{r} (1 - q^{-1} z_i^2) \prod_{i<j} (1 - q^{-1} z_i z_j)(1 - q^{-1} z_i z_j^{-1}).$$

Theorem 2 *We have*

$$D(\boldsymbol{z}; q)\, \Xi(\lambda, \boldsymbol{z}) = \boldsymbol{z}^{\lambda+\rho_C}\, W(\lambda),$$

where $W(\lambda)$ is the Whittaker value defined in (7).

Before beginning the proof, we first require the following simple result. Let

$$P = \prod_{i=1}^{r} (1 - q^{-1/2} z_i).$$

Lemma 1 *If $f = f^+$ then*

$$\frac{(f|w^{-1})(\boldsymbol{z})}{f(w\boldsymbol{z})} = \frac{P(w\boldsymbol{z})}{P(\boldsymbol{z})}.$$

Proof If $p(w, z) = P(w\boldsymbol{z})/P(\boldsymbol{z})$ then p satisfies the cocycle condition $p(ww', z) = p(w, w'z)p(w', z)$. The left-hand side also satisfies the same cocycle relation so we are reduced to the case where w is a simple reflection, in which case it follows easily from the definition. $\square$

Proof of Theorem 2 The function $\Xi(\lambda, \boldsymbol{z})$ in (13) may be rewritten in the form

$$\frac{\boldsymbol{z}^{\lambda+\rho_C}}{\Delta_C} \sum_{w \in W} (-1)^{l(w)} (\boldsymbol{z}^{-\lambda-\rho_C}|w) = \frac{\boldsymbol{z}^{\lambda+\rho_C}}{\Delta_C P(\boldsymbol{z})} \mathcal{A}(P(\boldsymbol{z})\boldsymbol{z}^{-\lambda-\rho_C}). \tag{14}$$

where the latter equality follows by replacing w by w^{-1} and using the Lemma. Again, we've employed the notation for the alternator as in (6). Note in particular that for any rational function f, $\mathcal{A}(f(\boldsymbol{z})) = \mathcal{A}(f(w_0\boldsymbol{z}))$. Since $w_0\boldsymbol{z} = (z_1^{-1}, \cdots, z_r^{-1})$, the expression on the right-hand side of (14) equals

$$\frac{\boldsymbol{z}^{\lambda+\rho_C}}{\prod_{i=1}^{r}(1 - q^{-1/2} z_i)} \frac{1}{\Delta_C} \mathcal{A}\left(\boldsymbol{z}^{\lambda+\rho_C} \prod_{i=1}^{r} (1 - q^{-1/2} z_i^{-1}) \right).$$

Multiplying by $D(\boldsymbol{z})$ and simplifying, the statement follows. $\square$

BZL Patterns

Let w_0 be the long Weyl group element. Choose a decomposition reduced decomposition $w_0 = s_{\omega_1} \cdots s_{\omega_N}$ into a product of simple reflections where $1 \leqslant \omega_i \leqslant r$ (the rank). Let

$$\omega = (\omega_1, \cdots, \omega_N)$$

be the corresponding *reduced word* for w_0.

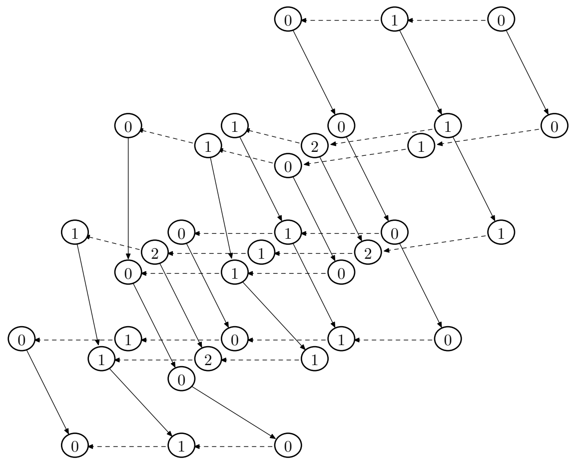
Let $\mathcal{B}_\lambda$ be the crystal of an irreducible finite-dimensional representation of highest weight λ for any Cartan type, and let W be the corresponding Weyl group. We will denote the Kashiwara (root) operators by e_i and f_i. The are maps $\mathcal{B} \longrightarrow \mathcal{B} \cup \{0\}$. There is a unique element $v_\lambda \in \mathcal{B}$ corresponding to the highest weight λ.

To each vertex $v \in \mathcal{B}_\lambda$ and each reduced word ω, we associate an integer sequence as follows. Let k_1 be the largest integer such that $e_{\omega_1}^{k_1}(v) \neq 0$. Then let k_2 be the largest integer such that $e_{\omega_2}^{k_2} e_{\omega_1}^{k_1}(v) \neq 0$, and so forth. Upon using all root operators in the order specified by the long word decomposition, $e_{\omega_N}^{k_N} \cdots e_{\omega_1}^{k_1}(v) = v_\lambda$, the vertex corresponding to the highest weight vector. The sequence $(k_1, \cdots, k_N)$ determines v, and can be arrayed in a pattern to give a convenient way of parametrizing elements of the crystal. These patterns were studied by Littelmann [18] and by Berenstein and Zelevinsky [2]. We will refer to the sequence $(k_1, \cdots, k_N)$ as a *BZL string* or *BZL pattern* and write $(k_1, \cdots, k_N) = \mathrm{BZL}_\omega(v)$. This construction applies equally well to any symmetrizable Kac-Moody group, but we focus entirely on type B root systems and their crystal graphs in this section.

Theorem 3 *Let $\mathcal{B}$ be a crystal graph of type B. There exists a unique function σ on $\mathcal{B}$ taking values in the nonnegative integers with the following properties. If v_λ is the highest weight vector then $\sigma(v_\lambda) = 0$. If $x, y \in \mathcal{B}$ and $f_i(x) = y$ with $i < r$, then $\sigma(x) = \sigma(y)$. If $e_r(x) = 0$, and $y = f_r^k(x)$, then*

$$\sigma(y) = \begin{cases} \sigma(x) & \text{if } k \text{ is even,} \\ \sigma(x) + 1 & \text{if } k \text{ is odd.} \end{cases}$$

Let us illustrate this with an example.



This illustrates the crystal with highest weight $\lambda = (2, 1)$ for B_2. We draw $x \longrightarrow y$ with a solid arrow if $y = f_1(x)$, and with a dashed arrow if $y = f_2(x)$. The vertex in the upper right-hand corner is v_λ. The values of σ are shown for every element.

Proof We will give one definition of σ for each reduced decomposition

$$w_0 = s_{\omega_1} s_{\omega_2} \cdots s_{\omega_{r^2}}$$

of the long element. We will show that these definitions are all equivalent, then deduce the statement of the theorem. We start with the BZL string of $v \in \mathcal{B}$ corresponding to this word. Thus corresponding to the word

$$\omega = (\omega_1, \omega_2, \cdots, \omega_{r^2})$$

we produce the sequence $k_1, k_2, \cdots, k_{r^2}$ with each k_n defined by

$$e_{\omega_n}^{k_n} \cdots e_{\omega_1}^{k_1} v \neq 0, \qquad e_{\omega_n}^{k_n+1} \cdots e_{\omega_1}^{k_1} v = 0,$$

so that $e_{\omega_{r^2}}^{k_{r^2}} \cdots e_{\omega_1}^{k_1} v = v_\lambda$ is the highest weight element of the crystal base. Define

$$\sigma_\omega(v) = \sum_{\omega_j = r} \begin{cases} 1 & \text{if } k_j \text{ is odd,} \\ 0 & \text{if } k_j \text{ is even.} \end{cases} \tag{15}$$

We wish to assert that if ω and $\omega' = (\omega'_1, \cdots, \omega'_{r^2})$ are two reduced decompositions then $\sigma_\omega = \sigma_{\omega'}$.

The proof will involve a reduction to the rank two case, so let us first prove that the statement is true for crystals of type B_2. In this case, there are only two reduced words and we may assume that $\omega = \{1, 2, 1, 2\}$ and $\omega' = \{2, 1, 2, 1\}$. In this case, Littelmann (cf. Section 2 of [18]) proved that

$$\begin{aligned}
k'_1 &= \max(k_4, k_3 - k_2, k_2 - k_1), \\
k'_2 &= \max(k_3, k_1 - 2k_2 + 2k_3, k_1 + 2k_4), \\
k'_3 &= \min(k_2, 2k_1 - k_3 + k_4, k_4 + k_1), \\
k'_4 &= \min(k_1, 2k_2 - k_3, k_3 - 2k_4).
\end{aligned}$$

From this it follows easily that the number of odd elements of the set $\{k'_2, k'_4\}$ is the same as the number of odd elements of the set $\{k_1, k_3\}$, that is, $\sigma_\omega = \sigma_{\omega'}$.

We turn now to the proof that $\sigma_\omega = \sigma_{\omega'}$ for arbitrary rank r. Consider the equivalence relation on all reduced words representing w_0 generated by $\omega \sim \omega'$ if ω' is obtained from ω by replacing a string $\{l, m, l, m, \cdots\}$ of length equal to the order N of $s_l s_m$ in the Weyl group by the string $\{m, l, m, l, \cdots\}$ of the same length. By a theorem of Tits, any two reduced decompositions are equivalent under this relation. As a consequence, it is sufficient to show that $\sigma_\omega = \sigma_{\omega'}$ when ω' is obtained by replacing an occurrence of l, m, l by m, l, m (if $m = l + 1 < r$), or an occurrence of l, m, l, m by m, l, m, l (if $l = r - 1, m = r$) or an occurrence of l, m by m, l when $|l - m| > 1$.

Suppose that $i_t = l$, $i_{t+1} = m$, etc. are the elements of ω that are changed in ω'. The elements i_{t-1} and i'_{t-1} of ω and ω' preceding this string (if it is not initial) are not l nor m, and similarly for the elements following it. Let

$$v_h = e_{i_h}^{k_h} \cdots e_{i_1}^{k_1} v, \qquad e_{i_h}^{k_h+1} \cdots e_{i_1}^{k_1} v = 0, \qquad v'_h = e_{i'_h}^{k'_h} \cdots e_{i'_1}^{k'_1} v, \qquad e_{i'_h}^{k'_h+1} \cdots e_{i'_1}^{k'_1} v = 0,$$

so $v_0 = v'_0 = v$ and $v_{r^2} = v'_{r^2} = v_\lambda$. We will argue that the sequences $(v_0, v_2, \cdots, v_{r^2})$ and $(v'_0, v'_2, \cdots, v'_{r^2})$ and the sequences $(k_1, k_2, \cdots, k_{r^2})$ and $(k'_1, k'_2, \cdots, k'_{r^2})$ are identical, except at indices t through $t + N - 2$ where ω and ω' differ.

To see this, remove all edges of the crystal graph except those labeled l and m, which produces a crystal graph $\mathcal{B}'$ of type A_2, B_2, $A_1 \times A_1$ or $A_1 \times B_1$. Clearly $v_{t-1} = v'_{t-1}$ since ω and ω' agree up to this point. Let $\mathcal{B}''$ be the connected component of $\mathcal{B}'$ containing this. Then v_{t+N-1} is the highest weight vector in $\mathcal{B}''$ and so is v'_{t+N-1}. It is now clear that the portion of the BZL pattern which lies within this crystal is the only part of $k_1, \cdots, k_{r^2}$ which is different from $k'_1, \cdots, k'_{r^2}$, and we have only to show that the number of k_i within this subpattern with $\omega_i = r$ such that k_i is odd is the same as for the k'_i. That is, we have reduced to the rank two case. If $\mathcal{B}'$ is of type

B_2 we have proven this, and the other three cases are trivial, since an A_2 or $A_1 \times A_1$ crystal has no edges of type r, while an $A_1 \times B_1$ crystal is just a Cartesian product.

Now let $1 \leqslant i \leqslant r$. To verify the assertion that $\sigma(x) = \sigma(y)$ if $f_i(x) = y$, choose a word ω whose first element $\omega_1 = i$. If $(k_1, \cdots, k_{r^2}) = \mathrm{BLZ}_\omega(x)$ then $(k_1 + 1, k_2, \cdots, k_{r^2}) = \mathrm{BZL}_\omega(y)$. Since $\sigma(x)$ is the number of odd k_i with $\omega_i = r$, it is obvious that $\sigma(x) = y$. On the other hand, suppose that $e_r(x) = 0$. Choosing ω such that $\omega_1 = r$, we have $\mathrm{BZL}(x) = (k_1, \cdots, k_{r^2})$ with $k_1 = 0$ while $\mathrm{BZL}(f_r^k(x)) = (k, k_2, \cdots, k_{r^2})$ and so obviously $\sigma(f_r^k(x)) = \sigma(x)$ if k is odd and $\sigma(x)+1$ if k is even. $\square$

We recall that the Weyl group acts on the crystal: each simple reflection s_i acts by reversing the i-root strings. It is shown that this action gives rise to a well-defined action of W on the crystal in Littelmann [19].

Proposition 1 *If λ is integral, then the function σ is constant on W orbits of the crystal.*

Proof It is clear from the definition that reversing the i-root string through $v \in \mathcal{B}_\lambda$ does not change $\sigma(v)$ if $i < r$ since σ is constant on the root string in that case. If $i = r$, then the fact that λ is integral means that each root string has odd length, and therefore $\sigma(s_i(v)) = \sigma(v)$ in this case also, since $v - s_r(v) = k\alpha_r$ with k even. (Note that if λ is half-integral, the Weyl group action does not preserve σ.) $\square$

Conjecture 1 *Assume that λ is integral. Then*

$$\frac{1}{\Delta_C}\mathcal{A}\left(\boldsymbol{z}^{\lambda+\rho_C}\prod_{k=1}^{r}(1-q^{-1/2}\boldsymbol{z}_i^{-1})\right) = \sum_{v\in\mathcal{B}_\lambda}(-q^{1/2})^{\sigma(v)}\boldsymbol{z}^{\mathrm{wt}(v)}.$$

This expresses the metaplectic Whittaker function (except for its normalizing constant) as a sum over the crystal. As noted in a previous section, the left-hand side may be expanded as a polynomial in q whose coefficients are composed of irreducible characters. Hence, the above proposition may be viewed as partial evidence for the conjecture. It has also been verified numerically for many choices of λ.

Decorated BZL Patterns

Let us disregard the normalizing constant (9) for the time being, and consider (10) to be the value of the the p-adic Whittaker function at t_λ in $\widetilde{\mathrm{Sp}}_{2r}(F)$, where λ is a dominant weight for $\mathrm{Spin}_{2r+1}(\mathbb{C})$. Strictly speaking, this only makes sense if λ

is integral. However if λ is half-integral, it is probable that this scenario can be extended, taking t_λ in $\widetilde{\mathrm{GSp}}_{2r}(F)$. In any case, (10) is defined whether λ is integral or half-integral.

We saw in (12) that when $q^{1/2}$ is specialized to -1 the value of (10) becomes the character χ_λ^B of an irreducible representation of $\mathrm{Spin}_{2r+1}(\mathbb{C})$. We will reinterpret this fact in terms of crystals, showing that for any q, the expression (10) may be interpreted as a deformation of χ_λ^B. Indeed, we give a conjectural expression for the metaplectic Whittaker function evaluated at t_λ as a sum over vertices in the crystal $\mathcal{B}_\lambda$ by making use of certain decorated BZL patterns.

That is, we decorate the BZL string $k_1, \cdots, k_N$ by drawing boxes or circles around some of the entries according to the following rules. For the boxing rule, if

$$f_{\omega_i} e_{\omega_{i-1}}^{k_{i-1}} \cdots e_{\omega_1}^{k_1}(v) = 0$$

then we box k_i. Concretely, this means that the path from v to v_λ that goes through

$$v, e_{\omega_1}^{k_1}(v), e_{\omega_2}^{k_2} e_{\omega_1}^{k_1}, \cdots$$

includes the entire ω_i-string passing through the vertex $e_{\omega_i}^{k_i} \cdots e_{\omega_1}^{k_1}(v)$. In this sense, we roughly think of the value k_i as being as large as possible, and cannot be increased.

The circling rule may be regarded also very roughly as signifying that the value k_i is as small as possible, and cannot be decreased. To make this precise for type B_r and for one particularly nice reduced word ω, we take a closer look at the BZL patterns as treated by Littelmann [18].

We will use the Bourbaki ordering of the weights, so that the fundamental dominant weights are $\omega_1, \cdots, \omega_r$ with $\omega_1 = (1, 0, \cdots, 0)$ the highest weight of the standard representation, and $\omega_r = \left(\frac{1}{2}, \cdots, \frac{1}{2}\right)$ the highest weight of the spin representation. Then the reduced decomposition that we will use is

$$w_0 = s_r(s_{r-1}s_r s_{r-1})(s_{r-2}s_{r-1}s_r s_{r-1}s_{r-1}) \cdots (s_1 \cdots s_r \cdots s_1).$$

Thus $\omega = (r, r-1, r, r-1, r-2, r-1, r, r-1, r-2, \cdots)$ and $N = r^2$. An alternative indexing will sometimes be convenient, so we will write alternatively

$$\mathrm{BZL}(v) = (k_1, \cdots, k_{r^2}) = (k_{r,r}, k_{r-1,r-1}, k_{r-1,r}, k_{r-1,r+1}, \cdots).$$

Following Littelman, we put the entries into a triangular array, from bottom to top and left to right, thus

$$\left\{\begin{array}{ccccccc} k_{1,1} & \cdots & & k_{1,r} & k_{1,r+1} & \cdots & k_{1,2r-1} \\ & \ddots & & \vdots & & \cdot^{\cdot^{\cdot}} & \\ & & k_{r-1,r-1} & k_{r-1,r} & k_{r-1,r+1} & & \\ & & & k_{r,r} & & & \end{array}\right\} = \left\{\begin{array}{ccccccc} \ddots & & & \vdots & & & \cdot^{\cdot^{\cdot}} \\ & k_5 & k_6 & k_7 & k_8 & k_9 & \\ & & k_2 & k_3 & k_4 & & \\ & & & k_1 & & & \end{array}\right\}$$

Littelmann proved that the entries in each row satisfy the following inequalities (independent of the choice of highest weight λ):

$$2k_{i,i} \geqslant 2k_{i,j+1} \geqslant \cdots \geqslant 2k_{i,r-1} \geqslant k_{i,r} \geqslant 2k_{i,r+1} \geqslant \cdots \geqslant 2k_{i,2r-i} \geqslant 0$$

Note that every value is doubled except the middle one.

We circle the BZL string entry k_i if the corresponding lower bound inequality is an equality. Let us make this explicit in the case $r = 3$. In this case

$$\mathrm{BZL}(v) = (k_{3,3}, k_{2,2}, k_{2,3}, k_{2,4}, k_{1,1}, k_{1,2}, k_{1,3}, k_{1,4}, k_{1,5}) = (k_1, k_2, \cdots, k_9)$$

and the array is:

$$\left\{ \begin{array}{ccccc} k_{1,1} & k_{1,2} & k_{1,3} & k_{1,4} & k_{1,5} \\ & k_{2,2} & k_{2,3} & k_{2,4} & \\ & & k_{3,3} & & \end{array} \right\} = \left\{ \begin{array}{ccccc} k_5 & k_6 & k_7 & k_8 & k_9 \\ & k_2 & k_3 & k_4 & \\ & & k_1 & & \end{array} \right\}. \tag{16}$$

We have

$$k_{3,3} \geqslant 0$$

and if $k_{3,3} = 0$ we circle it. We have $2k_{2,2} \geqslant k_{2,3}$ and if this is an equality, we circle $k_{2,2}$. Similarly $k_{2,3} \geqslant 2k_{2,3}$ and if this is equality, we circle $k_{2,3}$.

We attach a simple root of the B_r root system to each column of the array, in this order:

$$\alpha_1, \cdots, \alpha_{r-1}, \alpha_r, \alpha_{r-1}, \cdots, \alpha_1.$$

The assignment is chosen so that a BZL string entry is in the column labeled by α_i if the corresponding element of the long word ω is i. Thus letting c_i be the sum of the i-th column, we have

$$\mathrm{wt}(v) = \lambda - (c_1 + c_{2r-1})\alpha_1 - (c_2 + c_{2r-1})\alpha_2 - \cdots - c_r\alpha_r.$$

Only α_r is a short root.

A Tokuyama function on BZL patterns

Let p be a prime element in the nonarchimedean local field F. Let $(\ ,\ )_n$ be the local n-th power Hilbert symbol. For any m and non-zero $c \in \mathfrak{o}$ we define the n-th order Gauss sum

$$g_t(m, c) = \sum_{\substack{x \bmod c \\ \gcd(x,c)=1}} \psi\left(\frac{mx}{c}\right) (x, c)_n^t.$$

We will only need these for $t = 1, 2$. For a nonnegative integer a, we also use the shorthand notations

$$g_t(a) = g_t(p^{a-1}, p^a), \qquad h_t(a) = g_t(p^a, p^a).$$

In the special case $n = 2$, then all these Gauss sums may be made explicit. The Gauss sum $g_1(1, p)$ is a square root of q, which we will denote $q^{1/2}$; by choosing the additive character ψ correctly we may arrange that it is the positive square root. Assuming $n = 2$, we then have:

$$g_1(a) = q^{a-\frac{1}{2}}, \qquad h_1(a) = \begin{cases} q^{a-1}(q-1) & \text{if } a \text{ is even,} \\ 0 & \text{otherwise,} \end{cases}$$

and

$$g_2(a) = -q^{a-1}, \qquad h_2(a) = q^{a-1}(q-1).$$

We now assume that n is even, and that $\mathcal{B}$ is crystal of type B_r. If $v \in \mathcal{B}$ we define

$$G(v) = \prod_{k \in \mathrm{BZL}(v)} \begin{cases} q^{-k} h_t(k) & \text{if } k \text{ is unboxed and uncircled,} \\ q^{-k} g_t(k) & \text{if } k \text{ is boxed but not circled,} \\ 1 & \text{if } k \text{ is circled but not boxed,} \\ 0 & \text{if } k \text{ is both boxed and circled.} \end{cases}$$

where the subscript $t = t(k)$ in the cases above is 1 if the root corresponding to k is α_r, and $t = 2$ otherwise. This means that $t = 1$ if k is in the middle column of the BZL array (16), and $t = 2$ otherwise. Note that these differ from the weights used in [5] in two ways:

- Due to the presence of both long and short roots we have two kinds of Gauss sums, indexed by t.
- The factor is multiplied by q^{-k} which ultimately simplifies the formulas.
- We have made our BZL patterns using the e_i instead of the f_i. This makes no real difference.

Now let λ be a dominant weight. Then we claim that $G(v)$ is a Tokuyama function for the metaplectic Whittaker function. More precisely:

Conjecture 2 *Assume that λ is integral. Then with $W(\lambda)$ as in (7), we have*

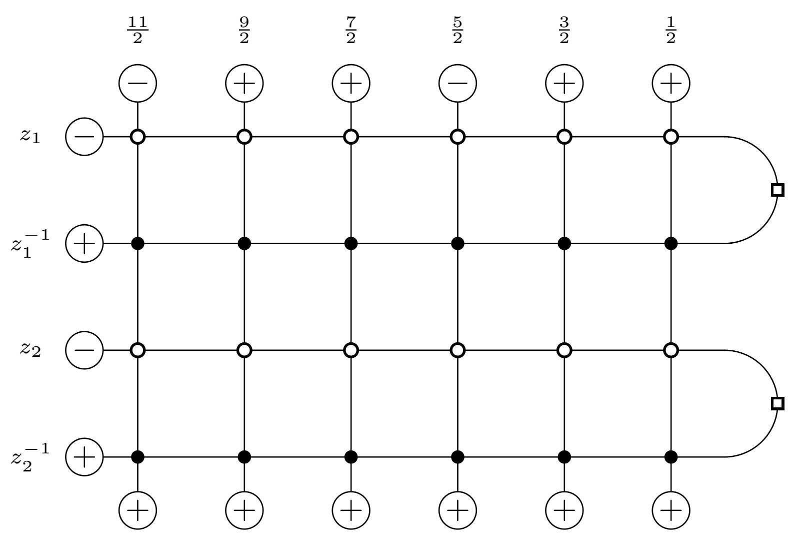
$$W(\lambda) = \sum_{v \in \mathcal{B}_{\lambda+\rho}} G(v) \boldsymbol{z}^{-\mathrm{wt}(v)}.$$

In order to make progress on this conjecture, we translate the problem to the realm of statistical lattice models. For more on the relationships between bases of highest weight representations and statistical lattice models, see [3] in this volume.

Ice Models

We will now give an alternative description of the Whittaker function as the partition function of a statistical system in the six-vertex model. The use of similar statistical models to study schur polynomials (i.e. p-adic Whittaker functions of type A with $n = 1$) was carried out by Brubaker, Bump and Friedberg [4]. The particular system we consider here is similar to the U-turn models used by Kuperberg [17] to enumerate classes of alternating sign matrices. A Tokuyama function for symplectic characters was discovered by Hamel and King [15] and understood in terms of the Yang-Baxter equation by Ivanov [16]. Despite these similarities, the particular model that we describe in this section is new and the first known instance of a Yang-Baxter equation for ice representing a metaplectic Whittaker function.

We consider a rectangular grid having $2r$ rows and the number of columns to be determined. The intersections of the rows and columns of the grid will be called *vertices*. The vertices in the odd numbered rows will be designated "Gamma ice" (labeled $\bullet$) and those in even numbered rows (labeled $\circ$) will be designated "Delta ice." Each pair of rows will be closed at the right edge by a "cap" containing a single vertex. Thus if $r = 2$ the array looks like:

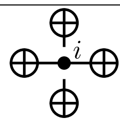
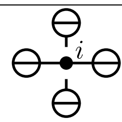
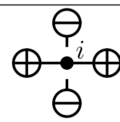
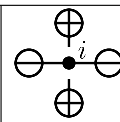
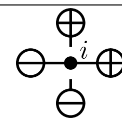
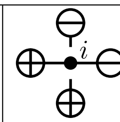
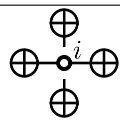
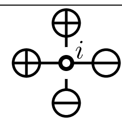
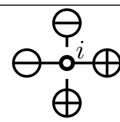
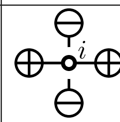
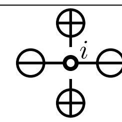
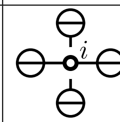


We have labeled the boundary edges by certain signs $\pm$. The interior edges will also be labeled with signs, but these signs will be variable, whereas the boundary edge signs are fixed and are part of the data describing the system.

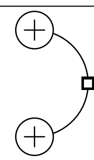
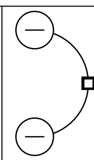
The boundary edge signs are to be assigned as follows. We put alternating signs $-, +, -, +, \cdots$ on the left edge, so that the rows of Delta ice begin with $-$ and the rows of Gamma ice begin with $+$. We put $+$ signs at the bottom of each column.

For the top, we label the columns with half integers beginning with $\frac{1}{2}$ at the right and increasing by 1 from right to left. Given a highest weight λ of type B, we put $-$ in the columns labeled from values in $\lambda + \rho_B$. Thus if $r = 2$ and $\lambda = (4, 2)$ then $\lambda + \rho_B = \left(\frac{11}{2}, \frac{5}{2}\right)$ and so we put $-$ in those columns, as indicated in the figure above. The remaining top edges are labeled $+$.

A *state* of the system is an assignment of spins $\pm$ to the remaining interior edges. For the Gamma and Delta vertices, the assignments must be taken from the following choices, known as *Boltzmann weights* in the language of statistical mechanics.

Γ vertex						
Boltzmann weight	1	z_i^{-1}	t	z_i^{-1}	$z_i^{-1}(t+1)$	1
Δ vertex						
Boltzmann weight	z_i	$z_i(t+1)$	1	$z_i t$	1	1

For the cap vertices, which we will label with a $\square$, the two adjacent edges must have the same sign, as follows.

$\square$ Cap Vertex		
Boltzmann weight	$-\sqrt{-t}z_i^{1/2}$	$z_i^{-1/2}$

For the moment, we may regard $t, z_1, \ldots, z_r$ as arbitrary parameters to be determined later. Given a highest weight λ of type B, we may fix boundary spins as above. Then an admissible state is one in which each vertex in the state has a Boltzmann weight taken from the above table. Let $\mathfrak{S}_\lambda = \mathfrak{S}_\lambda(z_1, \cdots, z_r, t)$ be the set of all states.

Given a state $S \in \mathfrak{S}_\lambda$ of the system, the Boltzmann weight $\mathrm{BW}(S)$ of the state is the product over all vertices of the weights of the vertex. The *partition function* $Z(\mathfrak{S})$ is the sum of the $\mathrm{BW}(S)$ over all states S.

As before, we let the Weyl group W of type B_r act on the parameters $z_1, \cdots, z_r$; it is generated by permutations of the z_i and the 2^r transformations $z_i \to z_i^{\pm 1}$.

Theorem 4 *The product*

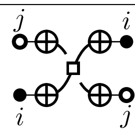
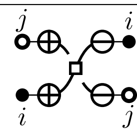
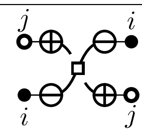
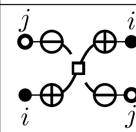
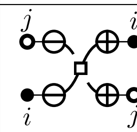
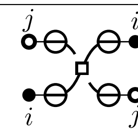
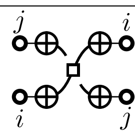
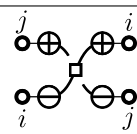
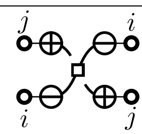
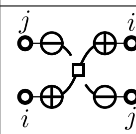
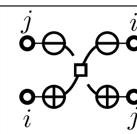
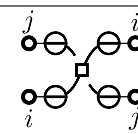
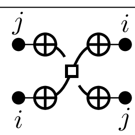
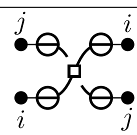
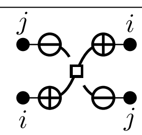
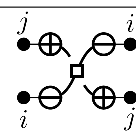
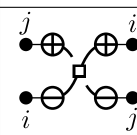
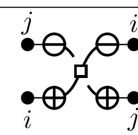
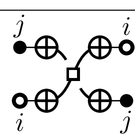
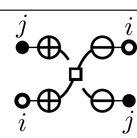
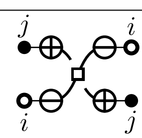
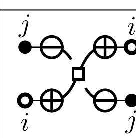
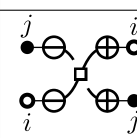
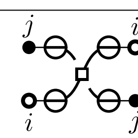
$$\boldsymbol{z}^{\rho_B}\prod_i(1-i\sqrt{t}z_i^{-1})\left[\prod_{i>j}(1+tz_iz_j)(1+tz_iz_j^{-1})\right]Z(\mathfrak{S}) \tag{17}$$

is invariant under the action of W*.*

The ideas of this proof are similar to those in [4] and [16], where the "caduceus" braid also appears.

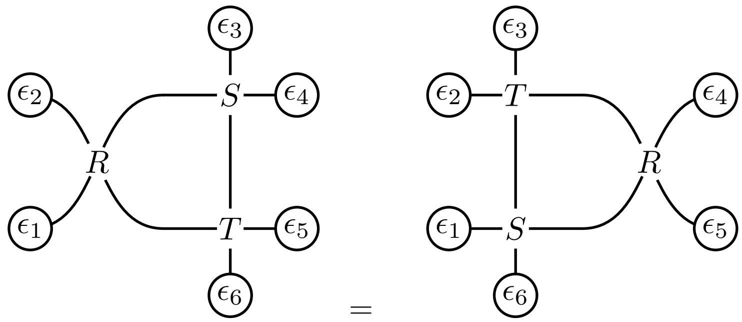
Proof We first show invariance under the simple reflections which interchange z_i and z_{i+1}.

We make use of the following types of vertices.

Type						
$\Gamma\Delta$						
Boltzmann weight	$t^2z_j-z_i^{-1}$	$(t+1)z_j$	$tz_j+z_i^{-1}$	$tz_j+z_i^{-1}$	$(t+1)z_i^{-1}$	$z_i^{-1}-z_j$
$\Delta\Delta$						
Boltzmann weight	tz_i+z_j	$z_j(t+1)$	tz_j-tz_i	z_i-z_j	$(t+1)z_i$	z_i+tz_j
$\Gamma\Gamma$						
Boltzmann weight	$tz_i^{-1}+z_j^{-1}$	$tz_j^{-1}+z_i^{-1}$	$tz_j^{-1}-tz_i^{-1}$	$z_i^{-1}-z_j^{-1}$	$(t+1)z_i^{-1}$	$(t+1)z_j^{-1}$
$\Delta\Gamma$						
Boltzmann weight	$z_i-z_j^{-1}$	$(t+1)z_i$	$tz_i+z_j^{-1}$	$tz_i+z_j^{-1}$	$(t+1)z_j^{-1}$	$-t^2z_i+z_j^{-1}$

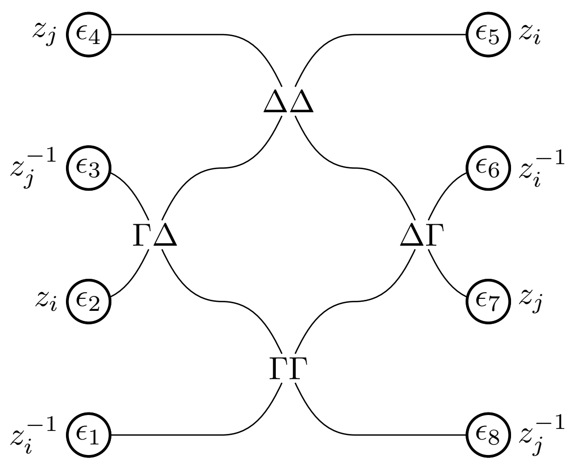
The results in [4] include the following "star-triangle relation" or Yang-Baxter equation. Given any pair $X,Y\in\{\Gamma,\Delta\}$, we may make three types of vertices: XY,

X and Y, each of whose Boltzmann weights are given by the above tables. Call these flavors of vertices R, S and T, respectively. Let $\varepsilon_1, \cdots, \varepsilon_6$ be six choices of sign $\pm$. Then the following two partition functions (each involving respective Boltzmann weights at three vertices) are equal.



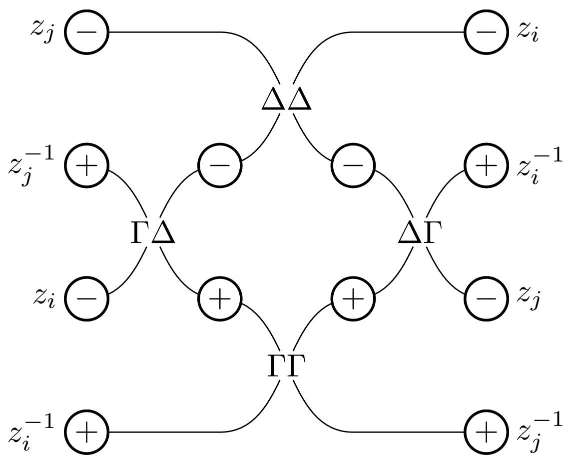
This means that (on each side of the equation) we sum over all assignments of signs to the 3 interior edges. The reversal of the spectral parameters, and of the order of the S and T vertices is indicated.

Now consider four rows of the system, which have (alternately) Δ, Γ, Δ, Γ vertices, with spectral parameters z_i, z_i^{-1}, z_j and z_j^{-1}. (So $j = i + 1$.) To the left of these four rows, we attach the following "caduceus" braid, which was employed in this context for type C by Ivanov [16].



We observe that there is only one legal configuration for this system which has

$(\varepsilon_5, \varepsilon_6, \varepsilon_7, \varepsilon_8) = (-, +, -, +)$. This configuration is:



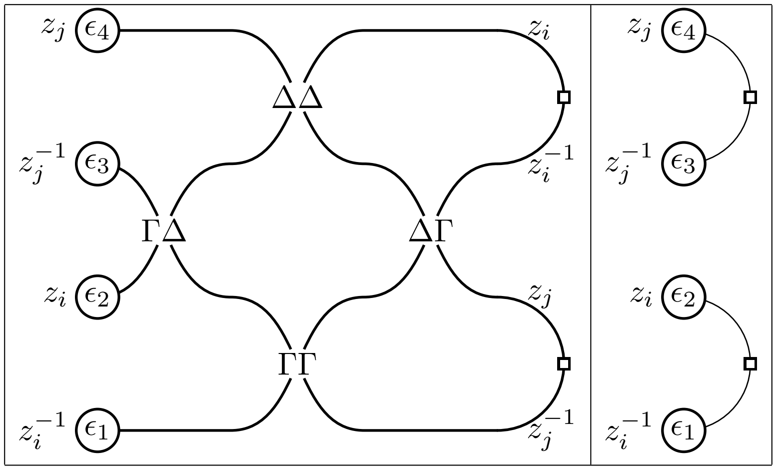
The partition function for this piece of ice is therefore just the product of the values at the four vertices, which can be read off from the above table:

$$(tz_j + z_i^{-1})(z_i + tz_j)(tz_i^{-1} + z_j^{-1})(tz_i + z_j^{-1}). \tag{18}$$

Hence we may attach the caduceus to the left of four rows in our original U-turn ice configuration and the resulting partition function multiplies the original partition function by this factor.

Using repeated application of the Yang-Baxter equation, the factor moves across the ice until it encounters the caps. In the process, the z_i and z_j spectral parameters are interchanged – effectively the two pairs of rows are switched. To move the caduceus past the caps along the right edge of the U-turn ice requires the following lemma.

Lemma 2 *Let $\varepsilon_1, \varepsilon_2, \varepsilon_3, \varepsilon_4 \in \{+, -\}$. Then the partition function of the system on* *the left in the following diagram*



equals

$$(tz_i + z_j^{-1})(tz_i + z_j)(tz_i^{-1} + z_j^{-1})(tz_j + z_i^{-1}). \tag{19}$$

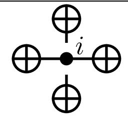
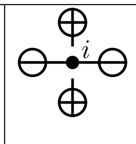
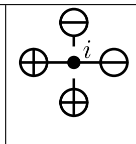
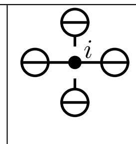
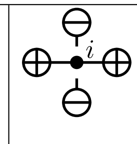
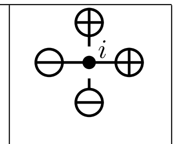
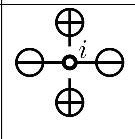
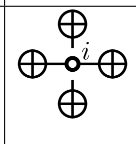
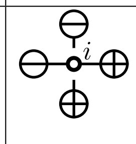
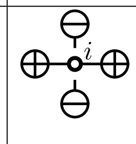
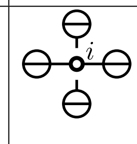
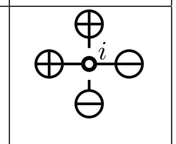
times the partition function of the system on the right.

This lemma allows us to pass the braid all the way through the U-turn ice, resulting in an interchange of parameters z_i and $z_j = z_{i+1}$. Thus our original partition function is related to that of U-turn ice with z_i and z_{i+1} swapped by the ratio of (18) to (19). This ratio equals

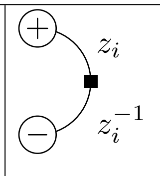
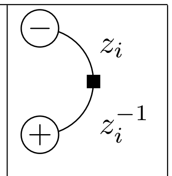
$$\frac{z_j + tz_i}{z_i + tz_j} = \frac{z_j}{z_i} \cdot \frac{1 - tz_i z_j^{-1}}{1 - tz_j z_i^{-1}} = \frac{\boldsymbol{z}^{s_i \rho_B}}{\boldsymbol{z}^{\rho_B}} \frac{1 - tz_i z_j^{-1}}{1 - tz_j z_i^{-1}},$$

which means that the product (17) is invariant under this interchange.

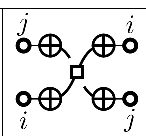
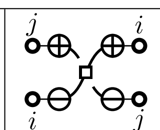
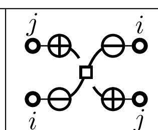
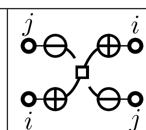
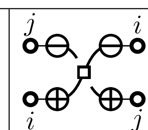
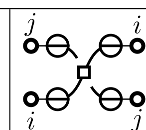
Now we consider the effect of the interchange $z_r \leftrightarrow z_r^{-1}$. For this, we begin by transforming the very bottom row of Γ vertices with spectral parameters z_r into Δ vertices with the spectral parameter z_r^{-1} by flipping the signs of all the horizontal edges in the row. Thus we will be using the following weights before and after the described change:

Γ vertex (before)						
Boltzmann weight	1	z_r^{-1}	1	z_r^{-1}	t	$z_r^{-1}(t+1)$
Δ vertex (after)						
Boltzmann weight	1	z_r^{-1}	1	$z_r^{-1}t$	1	$z_r^{-1}(t+1)$

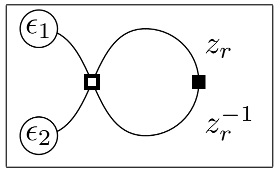
This change has no effect on the Boltzmann weights because of the boundary conditions. Indeed, only $+$ signs occur in the bottom edge spins, and therefore only the first three types of vertices in the table above occur. In order to compensate for the change, we must replace the cap vertices with the following modified ones, which we label by ■ instead of □:

■ Cap Vertex		
Boltzmann weight	$-\sqrt{-t}z_r^{1/2}$	$z_r^{-1/2}$

Now we attach a $\Delta\Delta$ vertex to the left, using the following Boltzmann weights:

$\Delta\Delta$						
Boltzmann weight	$tz_r + z_r^{-1}$	$z_r^{-1}(t+1)$	$tz_r^{-1} - tz_r$	$z_r - z_r^{-1}$	$(t+1)z_r$	$z_r + tz_r^{-1}$

As the signs in the bottom row have all been reversed, we are attaching the braid on the left to a pair of rows each beginning with $-$. There is only one such admissible $\Delta\Delta$ vertex – the last in the table above. Attaching this braid then multiplies our original partition function by $z_r + tz_r^{-1}$. We use the Yang-Baxter equation repeatedly to push this $\Delta\Delta$ vertex across the bottom two rows until it encounters the cap. Then we have the following configuration, referred to by Kuperberg [17] as the "fish equation":



It may be checked that the value of this configuration is

$$(1-\sqrt{-t}z_r)(1+\sqrt{-t}z_r^{-1})$$

times the value of the single ■ vertex. After this is substituted we may then repeat the flipping of all signs in the bottom row, turning the vertices in this row back into Γ vertices, but now with parameter z_r^{-1} changed to z_r.

Therefore $z_r + tz_r^{-1}$ times $Z(\mathfrak{S})$ equals $(1-\sqrt{-t}z_r)(1+\sqrt{-t}z_r^{-1})$ times the partition function with z_r replaced by its inverse. This implies that (17) is invariant under $z_r \to z_r^{-1}$. □

Based on the previous result, we may make the following conjecture.

Conjecture 3 *Take $t = -\frac{1}{q}$. Then $Z(\mathfrak{S})$ equals*

$$\frac{\boldsymbol{z}^{w(\rho_B)}}{\Delta_C} \prod_i (1+q^{-1/2} z_i) \left[\prod_{i<j} (1-q^{-1} z_i z_j)(1-q^{-1} z_i z_j^{-1}) \right] \mathcal{A} \left(\boldsymbol{z}^{\lambda+\rho_C} \prod_{k=1}^{r} (1-q^{-1/2} \boldsymbol{z}_i^{-1}) \right).$$

It follows from Theorem 4 that $Z(\mathfrak{S})$ is divisible by the product

$$\prod_i (1+q^{-1/2} z_i) \left[\prod_{i<j} (1-q^{-1} z_i z_j)(1-q^{-1} z_i z_j^{-1}) \right],$$

and the quotient is a polynomial in $q^{-1/2}$ and z_i, z_i^{-1} that is invariant under the Weyl group. Regarding both sides as polynomials in the arbitrary parameter $q^{-1/2}$, we may confirm the identity in the conjecture for the special values $q^{-1/2} = 0$ or 1. We are then able to prove the conjecture if $r \leqslant 3$ by bounding the size of the possible degree of $q^{1/2}$ in the resulting partition function and using the known pair of special values in $q^{-1/2}$.

Thus, this ice-type model conjecturally represents the Whittaker function. This conjecture implies Conjecture 2 using the bijection between states of U-turn ice with boundary corresponding to λ and vertices in the crystal $\mathcal{B}_\lambda$ of type B having the $G(v) \neq 0$. This bijection is implicit in [3] given the bijection between Gelfand-Tsetlin patterns and BZL patterns described by [18].

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